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A proton dose calculation code for treatment planning based on the pencil beam algorithm

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Abstract. We have developed a 3D dose calculation code for proton treatment planning. In order to achieve a fast and accurate calculation in inhomogeneous matter like a human body, the code employs the pencil beam algorithm where a broad therapeutic beam is decomposed into finite number of narrow beams before entering the inhomogeneous material and the dose contributions from the individual beams are summed up to reconstruct the dose distribution. Comparisons with detailed Monte Carlo calculations showed good agreements in dose distributions while keeping the manageable computational time for treatment planning.

1. Introduction

A proton beam has a sharper and better depth-dose characteristics than conventional X rays and therefore it should be in particular effective for treatments where a very high precision is required. As the proton treatment technique has been improving in dose delivery precision, it is becoming of more importance to accurately calculate the delivered dose distribution in a patient. The dose calculation program for this purpose must handle the beam modifying devices in the treatment system and a patient. The functions of the devices should be appropriately modeled and the internal structure of a patient should be given by the planning CT images.

The theory of physical interactions of a proton with matter has been well established and therefore the Monte Carlo technique should be able to generally simulate therapeutic proton beams. It is, however, not practical to apply the Monte Carlo method to clinical proton treatment planning since it requires too heavy computations to handle even on a latest computer system. Two important beam properties involved in the dose calculation are the longitudinal depth-dose characteristics and the lateral beam spreading due to scattering. By assuming that the medium is laterally uniform, which is called the infinite slab layer approximation, the beam shape can be represented by a few parameters which can be semi-analytically calculated. However, this approximation is applicable only when the scattering

effect in inhomogeneous matter is negligible though that is not the case for a large field therapeutic beam in a patient. In other words, if the lateral beam size is small enough with respect to the inhomogeneity, the infinite slab layer approximation should be a fair approximation. By decomposing a broad therapeutic beam into subbeams narrow enough to apply the approximation before entering an inhomogeneous material, we can then calculate the dose contribution from each subbeam at any point to reconstruct the dose distribution, which is called the pencil beam algorithm. Since the scattering effect in a human body is typically a cm or less which is smaller than the typical organ size, the pencil beam algorithm is therefore a reasonable approach for a fast and accurate dose calculation required for treatment planning.

Efforts have been made to realize a fast and accurate 3D proton dose calculation based on the pencil beam algorithm by Petti (1992), Hong *et al.* (1996), and Carlsson *et al.* (1997), for example. Carlsson *et al.*'s work was focused on the behaviors of proton pencil beams in water to generate a basic data set to be used in treatment plan optimization for the beam scanning technique while the other two were intended to be used for treatment planning like this work is. A notable difference between the present work and Petti's or Hong *et al.*'s is found in algorithm as to how the pencil beams are generated and how the dose distribution is calculated at each of the given points of interest (POI's) in patient. In our "beam-oriented" algorithm, multiple subbeams are generated not with respect to POI's but rather independently with sufficient interval and number to cover the given irradiation field and then each subbeam is transported into the human body which is represented by a set of voxels given by the CT images. The dose contributions from a subbeam to the POI's are then calculated and added to reconstruct the dose distribution. On the other hand, in Petti's or Hong's "POI-oriented" algorithm, about 100 subbeams are generated around each POI to calculate the dose at that point and this procedure independently loops over all the POI's. The advantage of the beam-oriented algorithm is that it allows more detailed calculation in beam transport since the number of beams is minimal. The disadvantage is that each of the generated beams has to be associated with the POI's around the beam path out of the large number of POI's distributed in 3D space, which requires a smart algorithm to minimize the computations.

This work is intended to provide a reasonably fast and reasonably accurate dose calculation tool for clinical treatment planning for conventional irradiation systems, for which the wobbler-scatterer method is currently in our scope as a primary option for the beam spreading system. In this paper, we present the algorithm, the physical models, and the evaluated performance of the program in comparison with detailed Monte Carlo calculations.

2. Materials and Methods

2.1. Algorithm and modeling of the devices

The program employs a straightforward algorithm in such a way that a beam is tracked from upstream down toward the patient. The beam is defined as a bundle of particles and is characterized by beam direction, number of particles or fluence F , mean kinetic energy E , range straggling σ_R and three parameters to describe the lateral beam spread, namely, spatial variance V_{xx} , angular variance $V_{\theta\theta}$, and covariance $V_{x\theta}$, where the angular and the spatial spreads are approximated by Gaussian distribution functions. Those beam parameters are modified by physical interactions as later described in the text together with their definitions.

2.1.1. Range Modulator

A range modulating device such as a ridge filter is modeled as a variable thickness flat degrader. The modulation is simulated by superposing repeatedly calculated dose distributions with various degrader thickness and irradiation weight, which specify the device shape in the case of a ridge filter. The number of modulation steps may be minimized by considering the range straggling or by introducing some additional beam energy spread, in order to reduce calculation time.

2.1.2. Scatterers/degraders

The upper part of the irradiation system is assumed to be laterally uniform and the scatterers, the degraders, and air gaps in between are modeled as a 1-dimensional stack of materials through which we transport a single beam down until the beam meets the device which has some lateral structure. The scatterers/degraders have material parameters necessary to calculate physical interactions of the beam in matter, which are density ρ , average atomic-mass number ratio Z/A , mean excitation energy I , radiation length X_0 , and nuclear interaction length at high energy λ_I . Parameters ρ , Z/A , and I for various materials are given by ICRU (1984) and X_0 and λ_I are given by Barnett *et al.* (1996) or are calculated according to the instructions given in the references.

2.1.3. Subbeam generation and wobbler system

The beam wobbling is approximated by a simple angular bending in a circular direction at a fixed height and is analytically included in the subbeam generation part of the code. That is, at the plane where lateral inhomogeneity begins, we decompose the parent beam into finite number of subbeams with infinitesimal spatial spread. The plane is divided into up to about 100×100 mesh-like pixels and we define a subbeam at center of each pixel. The subbeams directly inherit the properties of beam energy E and range straggling σ_R from the parent beam and the other subbeam parameters (marked by a prime) are analytically given by (Tomura 1997);

$$F'(r) = F \frac{\delta_x \delta_y}{2\pi V_{xx}} \exp\left(-\frac{r^2 + r_W^2}{2V_{xx}}\right) I_0, \quad (1)$$

$$\theta'(r) = \frac{r V_{x\theta}}{V_{xx}} - \left(\frac{r_W V_{x\theta}}{V_{xx}} - \theta_W\right) \frac{I_1}{I_0}, \quad (2)$$

$$V_{\theta\theta}'(r) = V_{\theta\theta} - \frac{V_{x\theta}^2}{V_{xx}} + \left(\frac{r_W V_{x\theta}}{V_{xx}} - \theta_W\right)^2 \left(\frac{1}{2} + \frac{I_2}{2I_0} - \frac{I_1^2}{I_0^2}\right), \quad (3)$$

where the parameters without a prime are ones for the parent beam, $I_k = I_k(r r_W / V_{xx})$ is given by the k -th order modified Bessel function, $r = \sqrt{(x^2 + y^2)}$ is the distance to the parent beam axis from the pixel, δ_x and δ_y are subbeam intervals, θ_W and r_W are the wobbling angle and radius at the plane, respectively, and θ' is the subbeam angle measured from the parent beam axis to determine the subbeam direction. We define the spatial variance V_{xx}' and the covariance $V_{x\theta}'$ for the subbeam as;

$$V_{xx}' = \frac{\delta_x^2 + \delta_y^2}{24}, \quad (4)$$

$$V_{x\theta}' = \frac{V_{x\theta}}{V_{xx}} V_{xx}', \quad (5)$$

where V_{xx}' is to give the spread corresponding to the uniform distribution in the pixel and $V_{x\theta}$ is to preserve the beam divergence due to scattering since $V_{xx}/V_{x\theta}$ corresponds to the mean distance to the scatterers distributed in the beam line. The subbeams are then transported independently through materials downward.

2.1.4. Compensator

Range compensators or boluses may be put in the beam line which are modeled as sets of mesh-like pixels with material thickness given to each pixel. For each of the subbeams, the compensator just behaves as a flat degrader with thickness defined at the pixel though which the beam passes, because of the infinite slab layer approximation.

2.1.5. Collimator

A beam collimator is defined by mesh-like pixels in a plane to represent the 2-dimensional collimator shape. A beam which hits a collimator pixel with "open" state will pass through the collimator and otherwise the beam is killed immediately. In this model, we ignore the thickness of the collimator and hence also the particles scattered by the edges of the collimator. Only the final collimator can be included while the other upstream ones are ignored. In a case where subbeams are generated before the collimator due to presence of a compensator upstream of the collimator, for example, a special corrective action to the beam spatial spread is applied. The parameter to give the beam size V_{xx}' is kept minimized until the beam comes down to the collimator plane. This is to enable the sharp cut-off at the collimator edges in order to correctly reproduce lateral penumbra, though the effect is very small in practice since the compensator is usually located very close to the collimator. In a usual case where the collimator is the first device with lateral structure, the subbeam pixels and the collimator pixels may be identically defined in the same plane.

2.1.6. Patient

A patient is modeled as a set of voxels of variable density water. The density of each voxel is derived from planning CT data so as to give the correct stopping power for therapeutic protons. The errors in scattering effect and nuclear effect due to this approximation are expected to be negligibly small in a typical treatment (Matsufuji 1997). Averaging of and/or interpolation from the CT data may be necessary to define voxels with the optimum size which may be a few mm in all three directions for a conventional large field irradiation. For simplicity of the code, a POI is defined at the center of each voxel, where dose value is calculated to construct 3D dose distribution. The dose value at a POI located at depth z along the beam and lateral distance r from the beam axis, $D(z,r)$, is then given by a simple Gaussian spreading of the linear energy transfer (LET) L which is the mean absorbed energy per unit length along the beam, as;

$$D(z,r) = L(z) \frac{1}{2\pi V_{xx}(z)} \exp\left(-\frac{r^2}{2V_{xx}(z)}\right). \quad (6)$$

The derivation of L is explained in the next subsection. Since we are not usually interested in fine structure of dose distribution within a voxel, we apply a limitation to V_{xx} by;

$$V_{xx} \geq \frac{w_x^2(1-u_x^2) + w_y^2(1-u_y^2) + w_z^2(1-u_z^2)}{24}, \quad (7)$$

where w_x , w_y , and w_z are the voxel widths and unit vector (u_x, u_y, u_z) represents the beam direction in the patient coordinate system. The right hand side term replaces V_{xx} when V_{xx} is smaller than that to calculate smeared dose instead of the point dose at the POI.

Since parameters L and V_{xx} are calculated sparsely at the transport steps, we apply the interpolation technique to obtain the value at the arbitrary depth. The dose at every POI is ideally calculated by summing up contributions from all the beams though it will be a very time-consuming computation to take all the combinations of the POI's and the beams. By ignoring the tail part of the Gaussian distribution beyond 3 standard deviations, we can restrict the region to look for the POI's to be associated to each beam and hence the combinations have been minimized. In the case of example 2 presented later in the text, which may be a typical example for treatment planning, about the same amount of CPU time was spent in the dose distributing part and in the beam transport part while the remaining overhead was negligible.

2.2. Modeling of physical interactions

2.2.1. Ionization energy loss

The mean energy loss rate per unit path length, or the stopping power S , for charged particles in matter due to ionization processes is given by the Bethe-Bloch formula (ICRU 1993). In a simplified form with such an approximation that the proton mass is much heavier than the electron mass and ignoring all the higher order corrections, we can write the formula as;

$$S = -\frac{dE}{dz} = 0.3071\rho \frac{Z}{A} \frac{1}{\beta^2} \left\{ \ln \left(\frac{1.022\beta^2}{I(1-\beta^2)} \right) - \beta^2 \right\}, \quad (8)$$

where β is the proton velocity in unit of speed of light. As to the units, MeV for energy, cm for length, and g/cm^3 for density are used unless noted, throughout this report. The higher order correction terms are usually very small and even much smaller than the error coming from the uncertainty of parameter I and hence are ignored. The proton range in water calculated by this formulation agrees with the fully corrected proton range table in the reference within 0.15 mm for incident kinetic energy everywhere up to 250 MeV. The relation between the proton kinetic energy E and the depth z is given by integrating this equation. In practice, the integration is done only once for water as the medium and values are stored in a table to associate kinetic energy E and water equivalent depth ζ which is defined by;

$$\zeta(z) = \int_0^z \frac{S(z')}{S_w(z')} dz', \quad (9)$$

where S_w is the stopping power for protons virtually passing through water with the given velocity. By numerically calculating ζ along the beam path, we can therefore obtain E at any depth by interpolation of the E - ζ table.

Energy spread due to fluctuation of the energy loss processes in thick material is well approximated by a Gaussian distribution (Fano 1964). Ignoring all the small higher order terms, the standard deviation of the Gaussian energy spread is given by;

$$\sigma_E^2 = \int 0.1569 \rho \frac{Z}{A} \frac{1 - \beta^2}{1 - \beta^2} \frac{1}{2} dz. \quad (10)$$

In practice, the integration is numerically performed step by step as depth and we convert the energy spread at the fixed position into the position spread in water equivalent depth at the fixed energy or so-called range straggling as;

$$\Delta \sigma_R^2 = \frac{\Delta \sigma_E^2}{S_w^2} = 0.1569 \rho \frac{Z}{A} \frac{1 - \beta^2}{1 - \beta^2} \frac{1}{2} \frac{\Delta z}{S_w^2}, \quad (11)$$

where prefix Δ represents the increase of the variable in the integration step.

The linear energy transfer from a proton L_p is then given by a convolution of the stopping power and the Gaussian function with standard deviation given by the range straggling, as;

$$L_p(z) = \int_0^\infty \frac{1}{\sqrt{2\pi}\sigma_R(z')} \exp\left(-\frac{\{\zeta(z) - \zeta(z')\}^2}{2\sigma_R(z')^2}\right) \frac{S(z')^2}{S_w(z')} dz'. \quad (12)$$

The histogram in figure 1 shows the derived L_p in water for one proton with 250 MeV incident kinetic energy while the open circles show the stopping power plotted at every mm before convolution of the straggling effect.

2.2.2. Inelastic Nuclear Interaction

A proton occasionally encounters the inelastic nuclear interaction with a nucleus in the medium, where a considerable part of the proton's kinetic energy is spent to destroy the nucleus and/or to generate secondary particles. In such cases, the charged secondary particles including the interacted proton are absorbed either immediately or in much shorter distance than the proton's original residual range. Neutrons and photons in the secondary particles usually escape far away from the interaction point and hence their total energy together with some spent for nuclear fission is seen as a missing energy. We included the effect of this phenomenon into LET calculation by a simple model.

For proton fluence reduction in water due to absorption, we employed the simplest fit to the experimental data by Lee *et al.* (1993), where the fluence linearly decreases as depth, and extended the model to other less important materials by scaling by nuclear interaction length at high energy. In a differential form, the proton fluence is given by;

$$\frac{dF}{d\zeta} = -F \frac{0.012}{1 + 0.012(R - \zeta)} \rho \frac{\lambda_{lw}}{\lambda_l}, \quad (13)$$

where R is the proton range in water for the given incident energy, $R - \zeta$ is hence the residual range, and λ_{lw} is the nuclear interaction length of water.

The dose contribution from secondary particles is formulated by the simplest model where a fixed energy escapes or is spent in the inelastic process and the rest of the energy is absorbed locally. The local energy deposition approximation is to assume that the proton and the charged secondary particles stop immediately after the interaction. The error from the approximation is expected to be large at material boundaries where density drastically changes such as body surface but relatively small in deeper region in body. In addition, the secondary particles should have less kinetic energy to move in the deeper region, which also justifies the approximation for therapeutic applications. We fitted the escaping energy to be

45 MeV to best reproduce the LET plots given in reference Carlsson *et al.* (1997) which were from model calculations based on the data from Seltzer (1993). The LET considering the inelastic nuclear interactions is then given by;

$$L = F \cdot L_p + \frac{-dF}{dz} (E - 45). \quad (14)$$

Figure 2 shows how the inelastic nuclear interaction is included in LET calculation in the code, where the solid, dashed, and dotted histograms represent the proton LET in water without the effect, one with the proton absorption effect, and one also with the contribution from the secondary particles, respectively, all normalized to one incident proton with kinetic energy of 250 MeV.

2.2.3. Multiple scattering

The multiple scattering of a charged particle in atomic Coulomb field is well represented by the Molière theory (Scott 1963) but the theory requires too complicated calculations to be used for this application. We therefore adopted the empirical approach to approximate the scattering angle by a Gaussian distribution with standard deviation given by the improved Highland formula, which was proposed by Lynch and Dahl (1991);

$$V_{\theta\theta}(z) = 13.6^2 \int_0^z \frac{dz'}{p^2 \beta^2 X_0} \left(1 + 0.038 \ln \int_0^z \frac{dz'}{\beta^2 X_0} \right)^2, \quad (15)$$

where p is the proton momentum in MeV/c and $\sqrt{V_{\theta\theta}}$ corresponds to the standard deviation of the distribution of the particle angle with respect to the beam axis projected to one lateral axis. These integrations are numerically performed in practice, reflecting the proton energy changes to p and β . The beam transport step Δz is determined as the step which changes the kinetic energy relatively by 10% but no smaller than 0.1 mm, under such a constraint that the step should not exceed the boundary of different materials. In order to correct the small changes of p and β in each step, we made the following corrections in the stepwise calculations;

$$\frac{1}{\beta^2} \rightarrow \left(1 + \frac{S\Delta z}{2E} \right) \frac{1}{\beta^2}, \quad (16)$$

$$\frac{1}{p^2 \beta^2} \rightarrow \left(1 + \frac{S\Delta z}{E} \right) \frac{1}{p^2 \beta^2}, \quad (17)$$

where β and p in right hand sides are calculated at the beginning of each step by their definitions from the proton mass and the given kinetic energy. These correction formulae are derived from the approximations such that S is inversely proportional to β^2 , that E is much smaller than the proton mass, and that $S\Delta z$ is much smaller than E . For example, the error in spatial spread of 250 MeV pencil beam in water due to the step size defined above amounted +0.3% compared to a very short step sized calculation, while, without the corrections, the error amounted -4.5%. Another small correction is applied to path length calculations since, due to angular spread, the average path length of the protons in a step is different from the step size which is the straight path along the beam. Except for the beam transport itself, the corrected step size is therefore used to all the calculations of the beam parameters, which is given by;

$$\Delta z \rightarrow (1 + V_{\theta\theta})\Delta z. \quad (18)$$

This detour correction is derived from the average of the detour factor $1/\cos\theta$ with the small angle approximation, where the particle angle θ is given by the Gaussian distribution with standard deviation $\sqrt{(2V_{\theta\theta})}$. Since the detour correction itself is in fact very small for a proton beam, we ignored the radial dependence of the detour correction.

The angular deflections are then converted to the spatial displacement as the particle moves. The Fermi-Eyges theory extended by Brahme *et al.* (1981) gives how to transport the beam to obtain the spatial spread. Defining the mean scattering power T in the beam transport step Δz by;

$$T = \frac{\Delta V_{\theta\theta}}{\Delta z}, \quad (19)$$

the covariance $V_{x\theta}$ is given by the numerical integration of;

$$\Delta V_{x\theta} = V_{\theta\theta}\Delta z + \frac{T}{2}(\Delta z)^2, \quad (20)$$

and the spatial variance V_{xx} is similarly given by the numerical integration of;

$$\Delta V_{xx} = 2V_{x\theta}\Delta z + V_{\theta\theta}(\Delta z)^2 + \frac{T}{3}(\Delta z)^3. \quad (21)$$

2.3. Comparisons with Monte Carlo calculations

2.3.1. Monte Carlo program GEANT

We evaluated the performance of the code by comparing with Monte Carlo calculations where behaviors of individual particles are tracked down as precisely as possible and hence the systematic errors coming from the various approximations are minimized. For this purpose, we used GEANT version 3.21 for availability, reliability, and applicability to general particle simulations, which was originally developed for Particle Physics experiments (CERN 1994). Except for the hadron cutoff energy in water which we set to 1 MeV, we ran GEANT with options given by default, which means that the Bethe-Bloch formula with Landau's fluctuation theory for ionization loss, GHEISHA package (Fesefeldt 1985) for hadronic interactions, and the Molière theory for the multiple scattering, are used. However, a part of the GEANT code was modified for our purpose. Since GEANT lacks the capability to correctly implement chemical compounds for materials, water is usually defined by a mixture of hydrogen atoms and oxygen atoms. This treatment is causing an error to the stopping power by up to several percent since parameter I in the Bethe-Bloch formula is a chemical property. Therefore we defined water as a non-compound material and made a temporary hardwired modification to subroutine GDRELX to use user-defined value for parameter I for water to reproduce the range data by ICRU (1993) for 150 MeV proton.

In the following examples, we will demonstrate the performance of our code in terms of precision and speed in comparison with GEANT. The speed was measured by the elapsed time of PA8000 160 MHz CPU on Hewlett Packard C160 UNIX workstation which we currently plan to use for the near future medical facilities. It is noted that GEANT is not a program optimized for treatment planning and therefore the speed comparisons in this report do not quantitatively evaluate the Monte Carlo approach for treatment planning in general.

2.3.2. Example 1: a pencil beam in water

We first tested the code in the simplest case. A 150 MeV proton beam was incident into water and we calculated 3D dose distribution with POI's distributed at every 0.5 mm in x and y directions and every 1 mm in z direction in the pencil beam code. In GEANT, a volume of water was divided into 1 mm thick slices along the beam direction and the energy depositions from 10000 incident protons and secondary particles were accumulated in a 2-dimensional histogram by depth and by radius, utilizing the axial symmetry of the geometry. The longitudinal division of the volume was necessary to correctly count the deposited energy in the depth bins.

The computational time amounted much less than a second for the pencil beam calculation and about 10 minutes for GEANT. Figure 3 shows the LET normalized to one proton, reconstructed from the 3D dose distribution, where the solid and the dashed histograms are from the pencil beam and the GEANT calculations, respectively. The disagreement amounted typically a few percent of the peak LET and, at peak, GEANT gives smaller LET by several percent which was the maximum deviation. The isodose plots in z - x plane are presented in figure 4, where isodose lines are plotted at every 10% of the Bragg peak dose. These isodose plots look quite similar but are not able to show the tails below the 10% line. Figure 5 show the lateral dose profile at the depth of Bragg peak. Due to the limitation of the Gaussian approximation with the tail cut off in the pencil beam code, the disagreement in the tail region is in a relative sense large but is very small in absolute value.

2.3.3. Example 2: a spreaded beam in patient

We also tested the code for a typical treatment case, which means that a pencil beam is fed into the irradiation system, where the beam is spreaded and collimated, and applied to a patient. A 150 MeV proton beam was incident into the irradiation system, for which we put a wobbler system to give 25 mrad bending at 289 cm, a 0.67 mm thick lead scatterer at 252 cm, a 2 cm thick lucite degrader at 200 cm, a 10 cm diameter ring collimator at 22.5 cm, and a patient with a flat body surface located 10 cm below the collimator's bottom surface. The patient was modeled by a water phantom, in which we put a cylinder of double density water ($\rho = 2 \text{ g/cm}^3$) in the region of depth in body 1 to 3 cm and radius 2 to 3 cm, as an example of inhomogeneity. In the pencil beam calculation, 64×64 subbeams with 1.5625 mm interval in two lateral directions were generated at the collimator plane and $\pi/4 \sim 80\%$ of them were transported into the $20 \times 20 \times 20 \text{ cm}^3$ patient volume consisting of $10^6 \text{ } 2 \times 2 \times 2 \text{ mm}^3$ voxels of water for each of which a POI was assigned at the center and the density was doubled if the POI was in the cylindrical region. On the other hand, in GEANT, one million of 150 MeV protons were fed into the irradiation system and the patient was modeled by water with a cylinder of double density water put in it, all divided along the beam axis at every 2 mm.

For the 3D dose calculation, the pencil beam code took 22 seconds while GEANT took about 1.5 hours. Figure 6 shows the isodose plots as functions of radius from the beam center and depth in body, where the hatched area corresponds to the double density cylindrical region. The overall shapes look quite similar though GEANT suffers from poor statistics at small radius. In order to see the details, we made slices of these plots along depth as in figure 7. The slices were made right on the double density region ($r = 2.5 \text{ cm}$) for the top plot, at boundary ($r = 3 \text{ cm}$) for the middle plot, and away from the inhomogeneity ($r = 4 \text{ cm}$) for the bottom plot. Though the systematic dose difference at Bragg peak between the pencil beam calculation and GEANT still remains in these plots as in the previous example, it is shown in the top plot that the pencil beam code well reproduces the scattered component coming around the double density region, which made the small second peak in the plot.

3. Discussion

The substantial LET disagreement from GEANT at Bragg peak may be originated from the simplified treatment of physical interactions, especially, in inelastic nuclear interactions. However, GEANT has not been fully certified to be correct for this kind of application and hence we made no effort for the pencil beam code to fit with GEANT results. It is noted that the LET could be given by an external look-up-table based on the measurements, which includes all nuclear effects without modeling, instead of the internally calculated LET table as described in the text. For each subbeam in the pencil beam code, the lack of the tails of lateral beam spread causes some dose enhancement in the central region of the spread. That is however less significant in real therapeutic situations where many similar subbeams contribute to the dose at a POI. The pencil beam method is to be viewed as an approximation of the Monte Carlo method by bundling the particles into much less number of narrow beams. In expense of faster calculation speed, it lacks capability to correctly handle secondary particles and large angle scatterings especially at the collimator edges. Also, the infinite slab layer approximation in the pencil beam algorithm can produce large errors when there are large inhomogeneities within the spread of each individual beam, in particular, when the inhomogeneities are coherently integrated along the beam path. However, the infinite slab layer approximation is indispensable for our speed requirement and for simplicity of the code. If we were to minimize those errors, a fast Monte Carlo method on a parallel computer system would be the best approach.

In example 2, the pencil beam calculation took 22 seconds for a 150 MeV unmodulated proton beam of 10 cm diameter field in a patient with 10^6 POI's. With the usual range modulation of order of 10 cm, this code requires about 10 to 20 calculations with different degrader thickness, which will increase calculation time by the same factor. We therefore conclude that this code can handle a typical treatment plan with manageable time of several minutes. We also plan to include an option to utilize interpolation technique for range modulation simulation to reduce the number of modulation steps. Though we haven't quantitatively optimized the control parameters in the code such as voxel widths and subbeam intervals, the optimum numbers for them may be both about or somewhat smaller than the beam spread. It is, however, meaningless to have too small voxels compared to the beam spread due to the limitation of the infinite slab layer approximation. In order to achieve fast 3D dose calculation, it is essential to minimize the number of voxels and the number of subbeams while keeping the required resolution since, roughly speaking, the number of dose computations is proportional to the product of number of subbeams and the average number of POI's for a subbeam to interact with. So far, we have implicitly assumed the one-to-one correspondence between voxels and POI's for simplicity. However, it is also possible to define POI's independently on the voxels. If we reduce the number of POI's, the calculation will be accelerated by the reduction factor. It may be true that the POI-oriented algorithm can have smaller number of dose computations when the number of POI's is greatly reduced. However, in such a situation that we have to know the full dose distribution in patient, which we believe is natural for treatment planning, the beam-oriented algorithm must be simpler and faster. It is noted that the performance in terms of speed greatly depends on the way the code was written in practice rather than the principles on which the code is based.

Though the code currently supports only the wobblers-scatterer method as a beam spreading system, it won't be difficult to implement more options such as the dual-ring double scatterer method and various scanning methods when necessary. Inclusion of biological effects, when the basic data and the practical models are established, will be the task to be done in the future.

4. Conclusion

We have developed a reasonably fast and reasonably accurate 3D proton dose calculation code based on the pencil beam algorithm with the infinite slab layer approximation. The beam modifying devices and the internal structure of a patient are well handled by the code, which has been verified by comparisons with the Monte Carlo calculations. The code is basically ready to be installed into treatment planning systems as a practical proton dose calculation tool.

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Figure Captions

Figure 1. LET from a 250 MeV proton in water. The stopping power for the mean kinetic energy at every mm is also plotted as open circles and the LET is shown as a histogram which includes the range straggling effect by Gaussian convolution.

Figure 2. LET normalized to one incident 250 MeV proton. The solid, the dashed, and the dotted histograms represent; without inelastic nuclear interactions, only with the proton absorption effect, and also with dose contribution from the secondary particles, respectively.

Figure 3. LET in water normalized to one incident 150 MeV proton, reconstructed from 3D dose distributions calculated by the pencil beam code (solid histogram) and by GEANT (dashed histogram).

Figure 4. Isodose plots in a cut-out plane. At left is shown the pencil beam calculation and at right is shown the GEANT calculation. The isodose lines are plotted at every 10% of the Bragg peak dose in each plot.

Figure 5. Lateral dose profile at Bragg peak in logarithmic scale. The solid and the dashed histograms show the pencil beam calculation and GEANT calculation, respectively.

Figure 6. Isodose plots in the modeled patient for the spreaded 150 MeV proton beam. At left is shown the pencil beam calculation and at right is shown the GEANT calculation. The hatched area in each plot represents the double density cylindrical region.

Figure 7. Longitudinal dose profile in the patient for the spreaded 150 MeV proton beam at three different radii; 2.5 cm (at top), 3.0 cm (at middle), and 4.0 cm (at bottom). The solid and dashed histograms represent the pencil beam calculation and GEANT calculation, respectively.

Figures

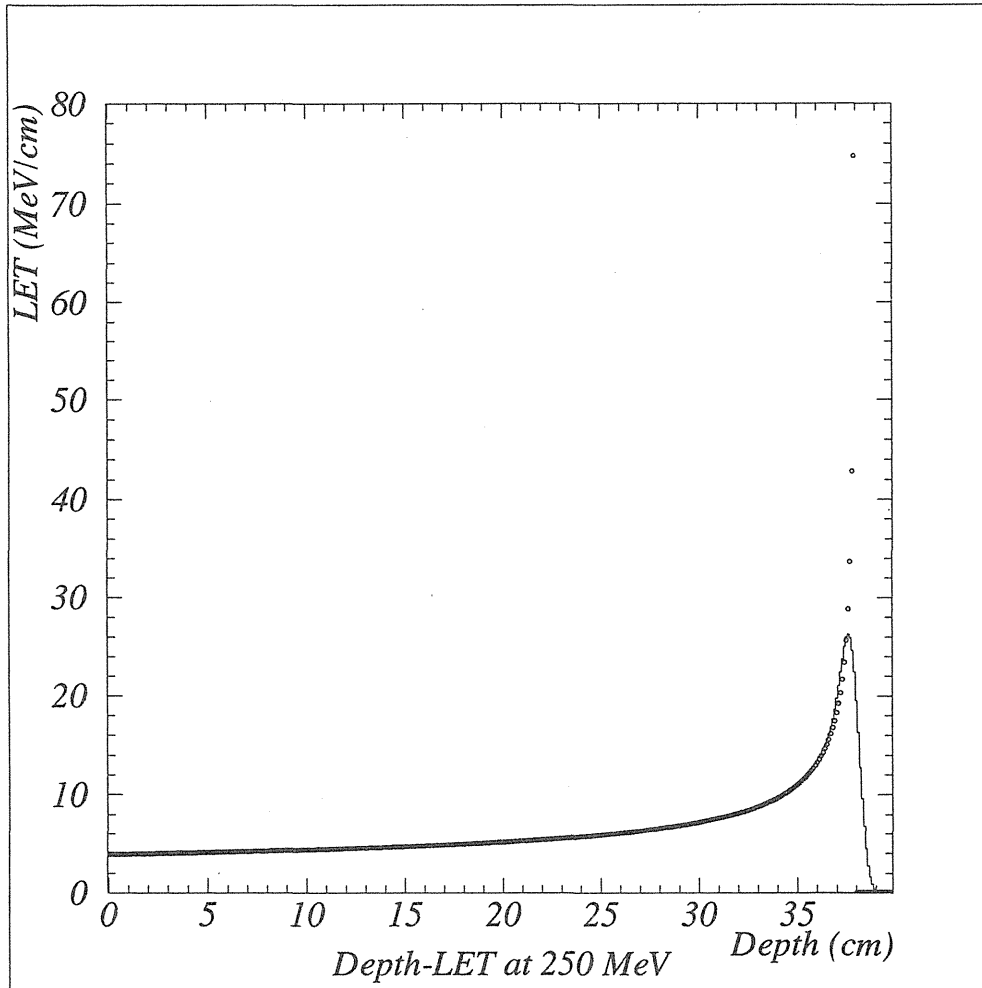


Figure 1. LET from a 250 MeV proton in water. The stopping power for the mean kinetic energy at every mm is also plotted as open circles and the LET is shown as a histogram which includes the range straggling effect by Gaussian convolution.

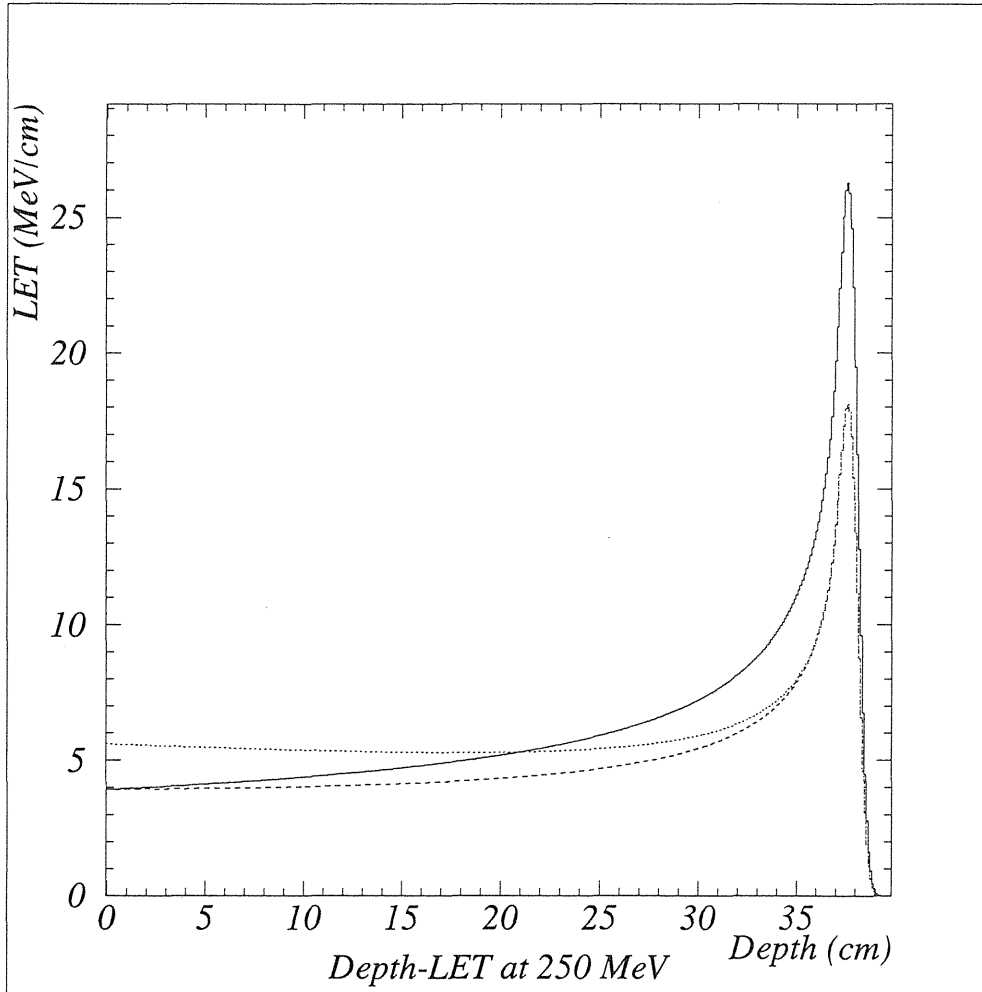


Figure 2. LET normalized to one incident 250 MeV proton. The solid, the dashed, and the dotted histograms represent; without inelastic nuclear interactions, only with the proton absorption effect, and also with dose contribution from the secondary particles, respectively.

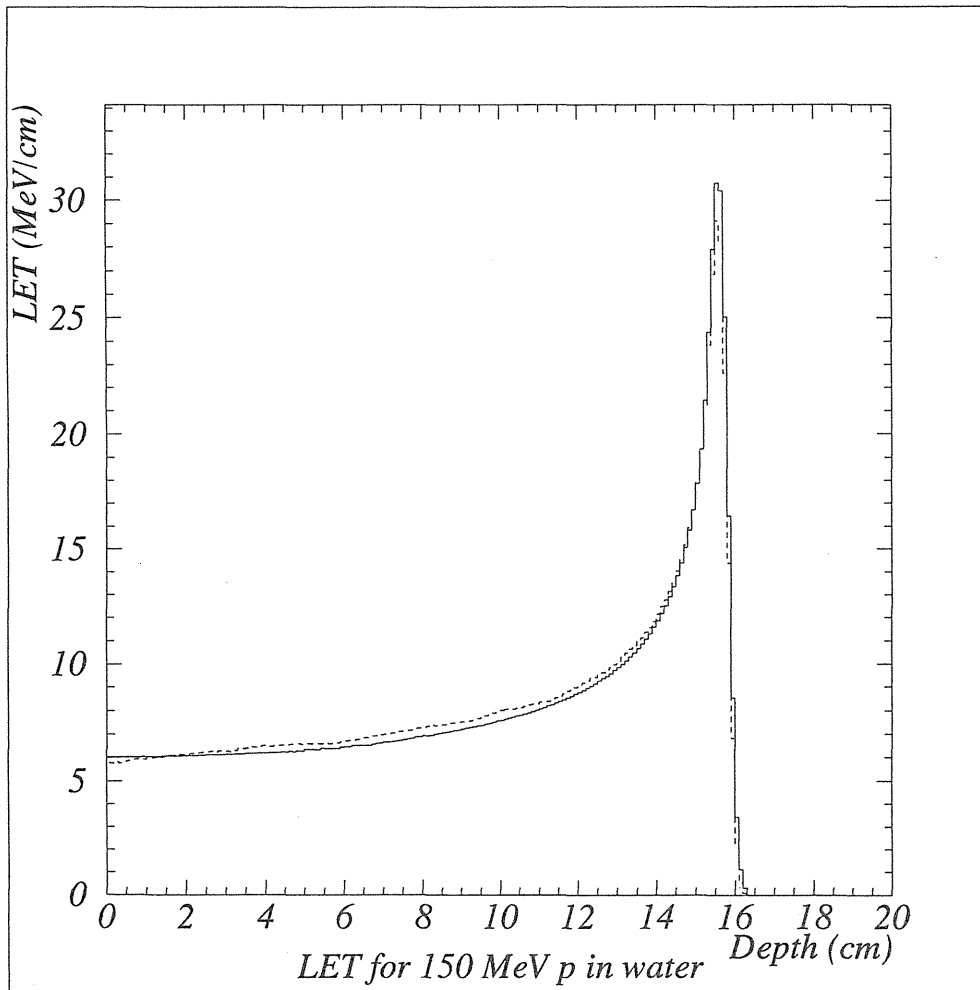


Figure 3. LET in water normalized to one incident 150 MeV proton, reconstructed from 3D dose distributions calculated by the pencil beam code (solid histogram) and by GEANT (dashed histogram).

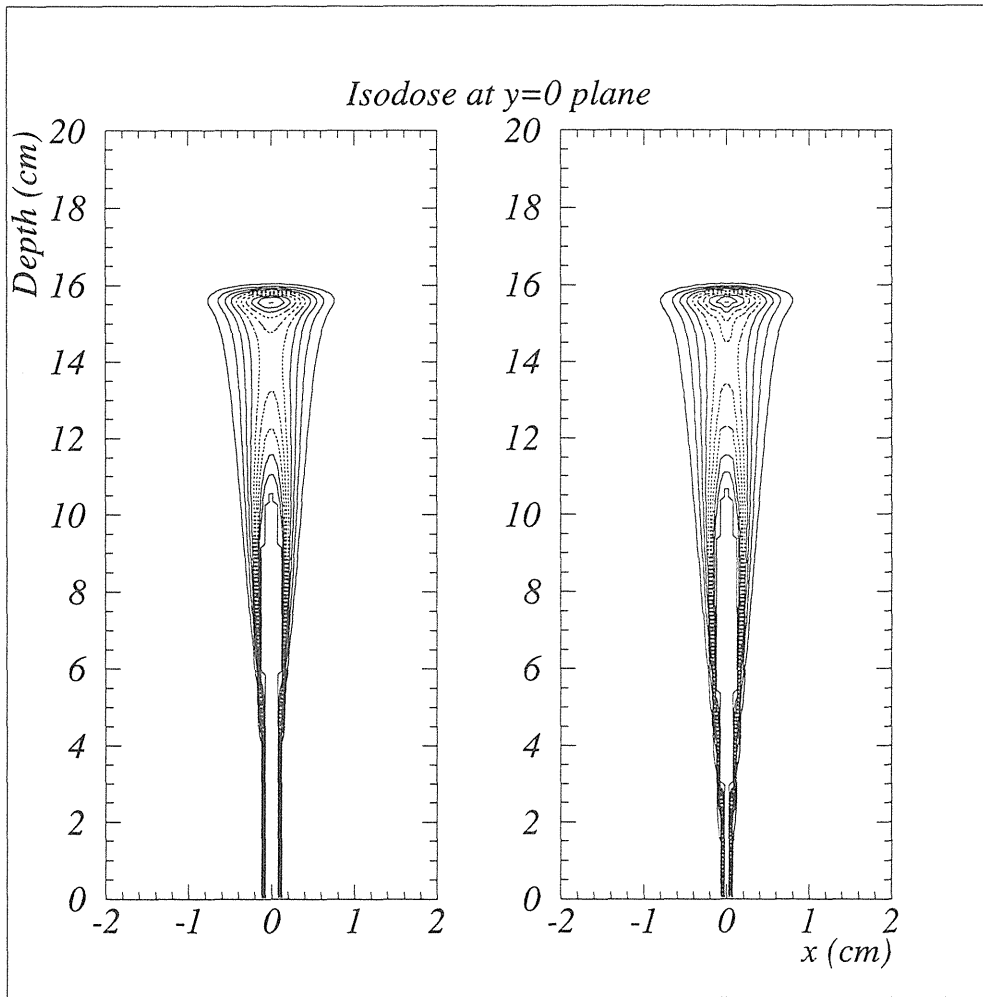


Figure 4. Isodose plots in a cut-out plane. At left is shown the pencil beam calculation and at right is shown the GEANT calculation. The isodose lines are plotted at every 10% of the Bragg peak dose in each plot.

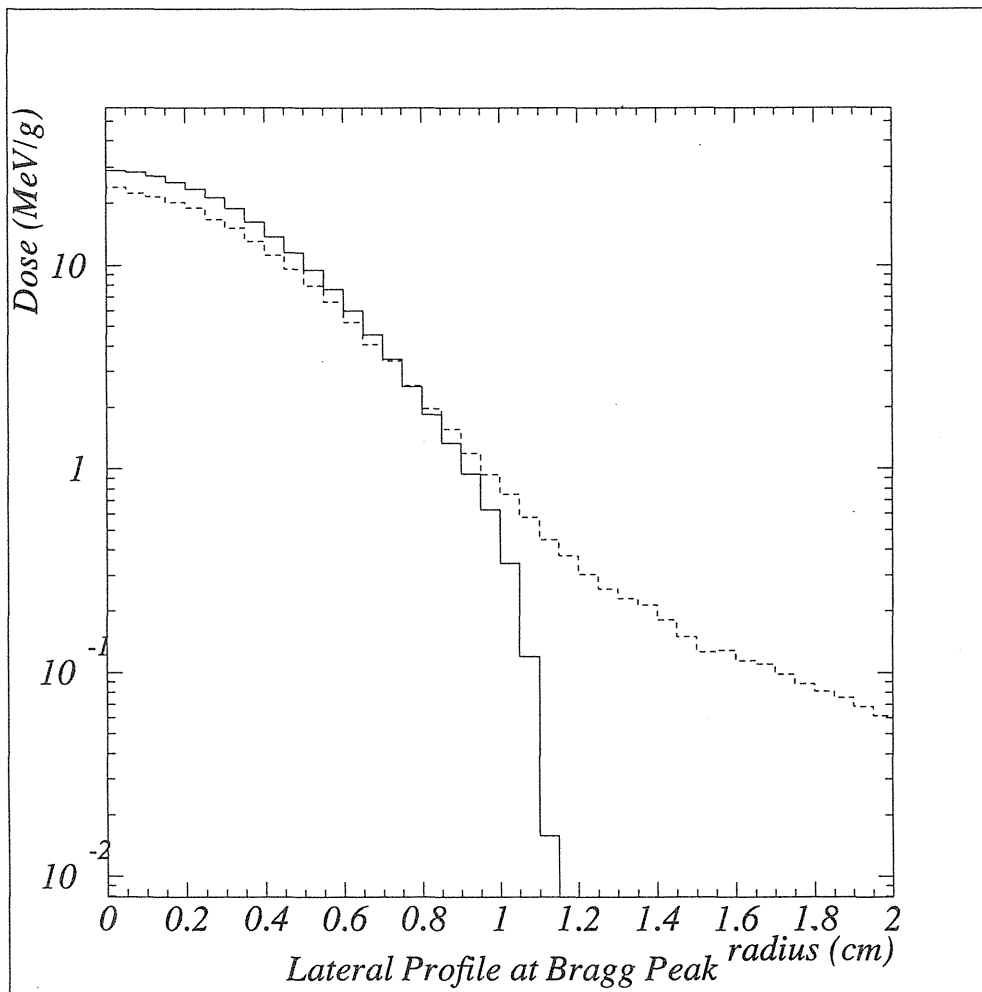


Figure 5. Lateral dose profile at Bragg peak in logarithmic scale. The solid and the dashed histograms show the pencil beam calculation and GEANT calculation, respectively.

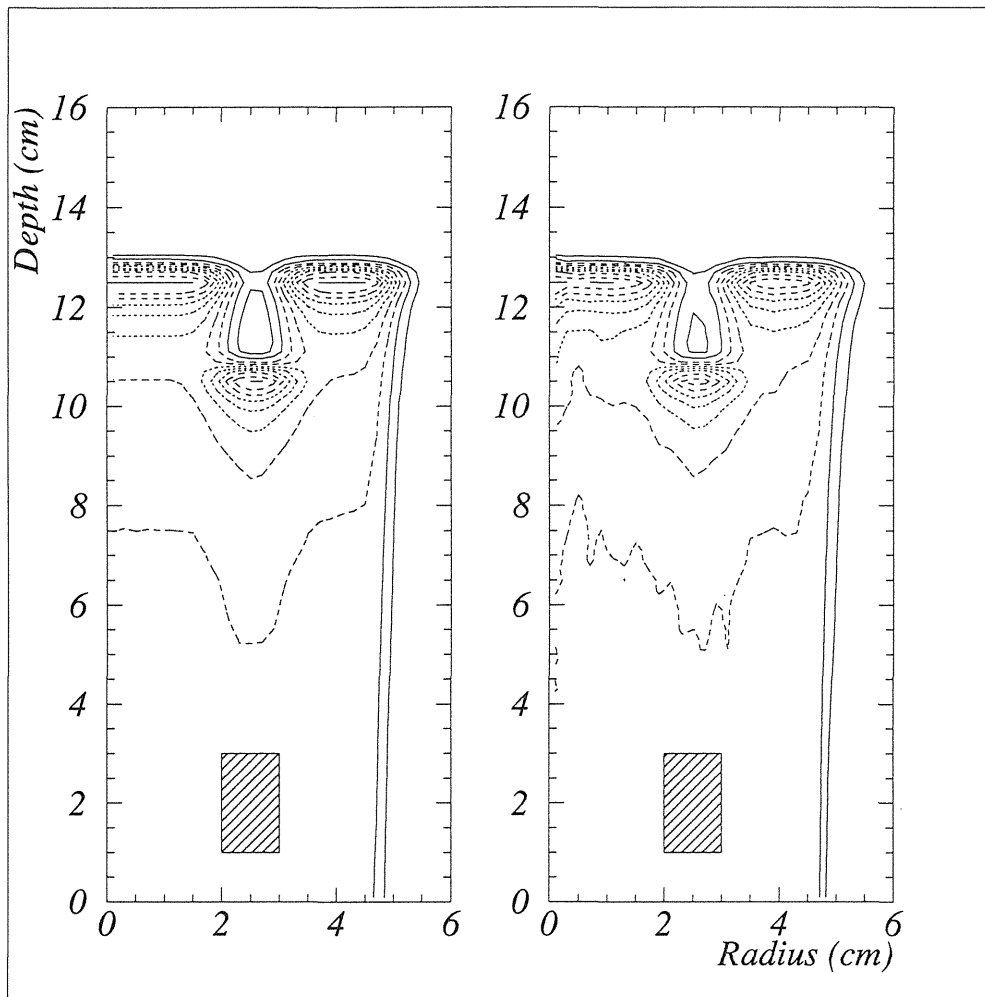


Figure 6. Isodose plots in the modeled patient for the spreaded 150 MeV proton beam. At left is shown the pencil beam calculation and at right is shown the GEANT calculation. The hatched area in each plot represents the double density cylindrical region.

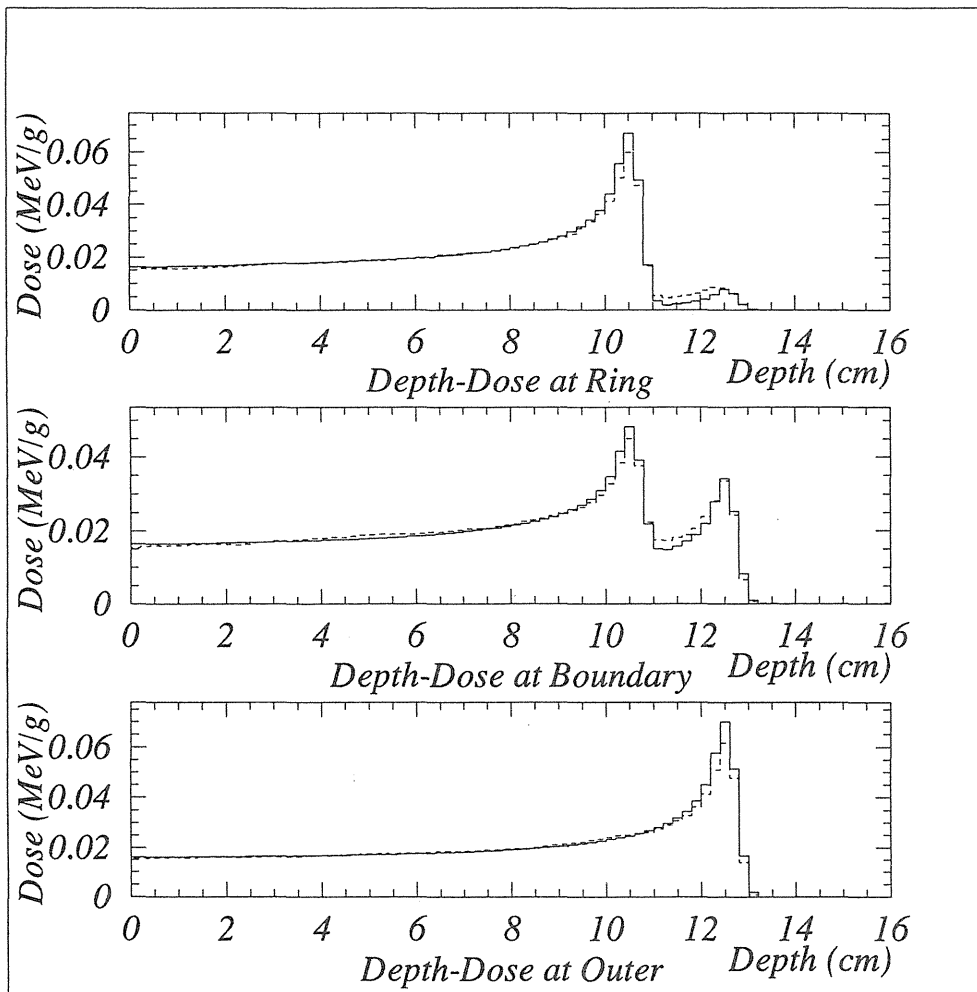


Figure 7. Longitudinal dose profile in the patient for the spreaded 150 MeV proton beam at three different radii; 2.5 cm (at top), 3.0 cm (at middle), and 4.0 cm (at bottom). The solid and dashed histograms represent the pencil beam calculation and GEANT calculation, respectively.

ペンシルビーム法陽子線量計算コード KDOSE マニュアル

1997年10月30日

三菱電機(株)通信機製作所 電子応用機器部 第一課 兼松 伸幸

(1997年2月～7月 放医研 重粒子治療センター 医用物理・工学研究部 研究生)

要旨

放医研・金井達明、外村浩美、松藤成弘、二見康之、山下晴男、兵庫県・赤城卓、東明男(敬称略)との共同研究において筆者が放医研研究生として開発した陽子線治療計画用ペンシルビーム法線量計算コードについてその内部を解説する。

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1.概要

1.1. 背景と目的

筆者は 1997 年 2 月より 7 月までの半年間放医研の研究生として重粒子治療センター医用重粒子物理・工学研究部第三研究室 金井達明室長らと共に陽子線治療計画用の線量計算コードの開発を行った。この計算コードは陽子線治療における治療計画作成のために人体内の線量計算を高速・高精度で行うことを目的とする。プログラムの名前は KDOSE とした。コード自身は 7 月中にほぼ完成し、概要および性能は日本医学物理学会の研究発表会(7 月 31 日茨城県つくば市)で発表し、ほぼ同じ内容で World Congress on Medical Physics and Biomedical Engineering(9 月 15 日フランス・ニース)においても発表した。また、同じ内容で英文の論文を作成中で、これは専門誌に投稿予定である。付録としてその原稿を添付する。このレポートは基本部分の完成した線量計算コードについてその内部構造を明らかにし、以後の応用や改良を容易にするためのものである。

1.2. 使用言語

FORTRAN77 の標準に準拠したが、拡張仕様の内、INCLUDE 文、SAVE 文、DO...ENDDO 文、DOWHILE...ENDDO 文、インラインコメント(!)は使用した。変数名、関数名、サブルーチン名、コメントブロック名はすべて英 6 文字以内とした。すべての関数及びサブルーチンに"IMPLICIT NONE"宣言を課し、全変数を宣言文で定義した。

1.3. 使用環境

上記の FORTRAN コードを正しくコンパイルできる FORTRAN コンパイラを必要とする。各社の UNIX ワークステーションならばすべて大丈夫と思われるが、これまで動作確認をしたのは HP(HPUX)、SGI(IRIX)、Sun(SunOS)である。線量計算コード自身は特殊なコンパイルオプションは必要としないが、誤差関数等の標準ライブラリは必要。HP では標準の誤差関数がなぜか正常に機能しなかったので CERN ライブラリ mathlib 中の誤差関数を使った。

1.4. アルゴリズム

いわゆるペンシルビームアルゴリズムを採用した。ただし、MGH 等で開発されたコードとは異なり、モンテカルロ法の簡略化というアプローチをとり、照射野全体にペンシルビーム(サブビーム)を走らせて各ビームから各計算点への線量の寄与を加算していく方法を採用した。その他、コードで採用した物理的なモデルや性能などの説明は添付の論文内で説明されているのでここでは省略し、このレポートではプログラムの内部構造のみを説明する。

1.5. 入出力

照射条件、患者データ等はコメントブロックの変数として計算に使われるが、現在のところ入出力に関しては FORTRAN のユーザールーチンを直接ユーザーが作成し、ここでコメント変数を操作することで行われる。従って、各ユーザーが各照射条件で実行ファイルを作成するか再コンパイルの必要のない入出力ルーチンを作成する必要がある。

1.6. プログラムの流れ

メインプログラム KPMAIN で最初に呼ばれる UKINIT でユーザーは照射パラメータ、患者データ等を定義する。レンジモジュレーションはディグレーダの厚さを変えて複数回線量計算を行い、結果を加算することで実現する。ビームを輸送は 2 段階で行われ、まず 1 本のビームをサブビーム生成面

まで輸送し、ここでのビーム形状から生成面の各ピクセルでサブビームを定義し、それぞれを輸送していく。ビーム輸送は定義された物体(volume)ごとに KSTRAK により行われる。患者内での線量分布は各サブビームからの各計算点での線量寄与を加算することにより求める。すべての計算を終えたら最後に UKLAST が呼ばれ、ここでユーザーは結果の解析あるいは他プログラムへの受け渡しを行う。

1.7. 使い方例

次節で説明する全ファイルを一つのディレクトリに置き、UNIX コマンド make を実行することにより、実行ファイル kdose ができる。これを実行することにより、ukinit.f で定義される照射条件による線量計算結果は uklast.f で指定されるようにヒストグラムパッケージ HBOOK を使って二次元分布がとられ、HBOOK ファイル wobblers.hbook に格納される。これを paw 等を使って画面上に表示させることができる。HBOOK 及び PAW は CERN ライブラリの一部である。サンプルプログラムでは HBOOK を線量分布表示のために使用しているだけで、そうしない場合は KDOSE を動かすのに CERN ライブラリは不要である。

1.8. バグ等

このコードをリリースする時点では既知のバグは存在しないが、数例に対してのみしかテストされていないので広く使われれば改修しなければならないバグが発見されるかもしれない。バグ報告やこのコードをより良いものに改良するアイデア、あるいは実際に改造して効果があったこと等あれば、以下に一報して頂けると助かります。

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2. ソースコードの詳細

2.1. メイクファイル

2.1.1. Makefile

```

SYSLIB = `cernlib mathlib packlib`

FFLAGS = -O -32
FFLAGS = -g -C -trapuv -32
FFLAGS = -g -C +ppu
FFLAGS = +O2 +ppu -G

INCLUDE = -I.
F77      = f77

all      : kdose

.SUFFIXES: .o .f

.f.o     :
$(F77) -c $(FFLAGS) $(INCLUDE) $<

kdose : rklet.o rkint1.o rkint2.o ¥
      dkaa.o kbmate.o kpmain.o ksptmt.o rkdxde.o rkintp.o ¥
      dkbesi.o kbmodu.o ksptst.o rkekin.o rkresr.o ¥
      rkenuc.o rksimp.o kbpai.o kstrak.o rkerf.o rkstop.o ¥
      kbbeam.o kbphys.o kswobb.o rkestr.o ikpvox.o ikpxyz.o ¥
      ksmeasr.o rkdfdr.o rkierf.o ¥
      ksinit.o ukinit.o uklast.o ksptdo.o
      $(F77) $(FFLAGS) $(INCLUDE) -o kdose ¥
      rklet.o rkint1.o rkint2.o ¥
      dkaa.o kbmate.o kpmain.o ksptmt.o rkdxde.o rkintp.o ¥
      dkbesi.o kbmodu.o ksptst.o rkekin.o rkresr.o ¥
      rkenuc.o rksimp.o kbpai.o kstrak.o rkerf.o rkstop.o ¥

```

CERN ライブラリとリンク。HP では誤差関数を含む mathlib も必要であったが他では必要なし。

IRIX 用コンパイルオプション(CERN ライブラリとリンクするため 32 ビットモードにしている)
 同、デバッグ用
 HPUX デバッグ用コンパイルオプション
 同、非デバッグ用
 ハードウェアや使用目的に応じて適当にオプションを指定する。

インクルードファイルの場所
 FORTRAN コンパイラ名

CERN ライブラリを使用するには、環境変数 CERN を CERN ライブラリを含むディレクトリツリーの元のディレクトリ名で定義し、環境変数 CERN_LVL はそのディレクトリの下位に位置しリリースバージョンに相当するディレクトリ名で定義しておく。さらにコマンドパスに \$CERN/\$CERN_LVL/bin を加えておく。このメイクファイル内で使用しているコマンド cernlib はそこに置かれたシェルスクリプトである。
 もしも UKLAST で HB00K を使用しない場合は CERN ライブラリに関係するコンパイルオプション (IRIX では -32, HPUX では +ppu) は削ってもよい。IRIX で 64 ビットモードあるいは N32 モードでコンパイルすると速度が格段に向上するはず。

<pre> kbbeam. o kbphys. o kswobb. o rkestr. o ikpvox. o ikpxyz. o ¥ ksnear. o rkdfdr. o rkierf. o ¥ ksnit. o ukinit. o uklast. o ksptdo. o \$(SYSLIB) </pre>	
--	--

2.2. インクルードファイル

2.2.1. *kcbeam. i*

<pre> *----- KCBEAM - Initial Beam -----* * REAL CHRGBM ! charge [e] REAL AMASBM ! mass [MeV/c^2] REAL FLUEBM ! fluence (number of particles) REAL EKINBM ! kinetic energy [MeV] REAL SEKBM ! spread of kinetic energy [MeV] REAL XYZBM(3) ! beam location REAL UVWBM(3) ! beam direction * COMMON/KCBEAM/CHRGBM, AMASBM, FLUEBM, EKINBM, SEKBM, XYZBM, UVWBM * *----- End of KCBEAM -----* </pre>	<p>照射系の受け取るビームの初期値を与えるためのコメントブロック</p> <p>電荷(陽子は+1) 質量(同 938MeV) 粒子数 運動エネルギー エネルギーの広がり(ガウシアン標準偏差) ビーム入射位置 入射方向</p> <p>KCBEAM あるいは UKINIT で定義される変数</p>
--	--

2.2.2. *kcbo/u. i*

<pre> *----- KCB0LU - Bolus -----* * INTEGER MBBO ! max number of boluses INTEGER MXBO ! max number of x division INTEGER MYBO ! max number of y division PARAMETER (MBBO=2, MXBO=256, MYBO=256) * INTEGER IVOLBO(MBBO) ! volume number INTEGER NXYBO(2, MBBO) ! number of x, y divisions REAL XY1BO(2, MBBO) ! lower left corner of first pixel [cm] REAL DXYBO(2, MBBO) ! pixel size [cm] REAL HIGHBO(MXBO, MYBO, MBBO) ! signed height at each pixel [cm] * COMMON/KCB0LU/IVOLBO, NXYBO, XY1BO, DXYBO, HIGHBO </pre>	<p>ボータスの形状データを与えるためのコメントブロック</p> <p>最大 MBBO 個のボータスが MXBO×MYBO の解像度で定義できる。</p> <p>/KCVOLU/でこのボータスに対応する物体の番号 ピクセル数 ボータスの角の位置 ピクセルの大きさ ボータスの各ピクセルの厚さ</p> <p>UKINIT で定義される変数</p>
--	--

```

*----- End of KCB0LU -----*
*-----*

```

2. 2. 3. kccoll. i

```

*----- KCCOLL - Collimator -----*
*
INTEGER MXCO      ! max number of x divisions
INTEGER MYCO      ! max number of y divisions
PARAMETER (MXCO=512, MYCO=512)
*
INTEGER IVOLCO    ! volume number
INTEGER NXYCO(2) ! number of x divisions
REAL XY1CO(2)    ! lower left corner of first pixel
REAL DXYCO(2)    ! pixel size
LOGICAL OPENCO(MXCO, MYCO) ! true if pixel is open
*
COMMON/KCCOLL/IVOLCO, NXYCO, XY1CO, DXYCO, OPENCO
*
*----- End of KCCOLL -----*

```

最終コリメータ形状を与えるためのコモンブロック
 最大 MXCO×MYCO の解像度のピクセルで定義する。
 /KVOLU/でこのコリメータに対応する物体の番号
 ピクセル数
 コリメータの角の位置
 ピクセルの大きさ
 各ピクセルが開いているかどうか
 UKINIT で定義される変数

2. 2. 4. kcdedo. i

```

*----- KCDED0 - depth-dose curve in water -----*
*
INTEGER MDDD ! max number of data points
PARAMETER (MDDD=256)
*
REAL EKINDD ! kinetic energy at depth=0.
INTEGER NDDD ! number of depth points
REAL DEPTDD(MDDD) ! depth in cmH2O
REAL DOSEDD(MDDD) ! measured LET in MeV/cm
*
COMMON/KCED0/EKINDD, NDDD, DEPTDD, DOSEDD
*
*----- End of KCDED0 -----*

```

実験値から深さ-線量特性を与えるためのコモンブロック
 最大データ点数
 深さ 0 でのビームの運動エネルギー
 実際のデータ点数
 深さ
 LET
 UKINIT で定義される変数

2. 2. 5. *kcmate. i*

-----	KOMATE - Material Data	-----	
*	INTEGER MAXMMT PARAMETER (MAXMMT=20)	! max number of materials	*
*	INTEGER IAIRMT INTEGER IWATMT CHARACTER*16 NAMEMT (MAXMMT) REAL ZOVAMT (MAXMMT) REAL EXTEMT (MAXMMT) REAL DENSMT (MAXMMT) REAL RADLMT (MAXMMT) REAL ABSLMT (MAXMMT)	! material number for air gaps ! material number for water ! material name ! effective Z/A ! mean excitation energy in MeV ! density in g/cm ³ ! radiation length in g/cm ² ! nuclear absorption length in g/cm ²	*
*	COMMON/KCMATE/IAIRMT, IWATMT, NAMEMT, ZOVAMT, EXTEMT, DENSMT, RADLMT &, ABSLMT		*
*	-----	End of KCMATE	-----*

物質データを与えるためのコモンブロック

最大 MAXMMT 個の物質が定義できる。

空気に対応する物質番号
水に対応する物質番号

各物質の名前

Bethe-Bloch の式のパラメータ

同上

物質の密度

クーロン散乱計算のための輻射長

高エネルギーでの核反応の特性長

KBMATE あるいは UKINIT で定義される変数

物質データは IORU37 (IORU) 及び Review of Particle Properties (Particle Data Group) による。

2. 2. 6. *kcmon. i*

-----	KOMCON - Mathematical Constants	-----	
*	REAL*8 M_E REAL*8 M_LOG2E REAL*8 M_LOG10E REAL*8 M_LN2 REAL*8 M_LN10 REAL*8 M_PI REAL*8 M_SQRT2 REAL*8 M_SQRTPI	! exp(1) ! log ₂ (e) ! log ₁₀ (e) ! ln(2) ! ln(10) ! pi ! sqrt(2) ! sqrt(pi)	*
*	PARAMETER (M_E = 2.7182818284590452354) PARAMETER (M_LOG2E = 1.4426950408889634074) PARAMETER (M_LOG10E = 0.43429448190325182765)		*

数学的な定数の定義

```

PARAMETER (M_LN2      = 0. 69314718055994530942)
PARAMETER (M_LN10     = 2. 30258509299404568402)
PARAMETER (M_PI       = 3. 14159265358979323846)
PARAMETER (M_SQRT2    = 1. 41421356237309504880)
PARAMETER (M_SQRTPI   = 1. 77245385090551602729)
*
*----- End of KCMCON -----*

```

2. 2. 7. kcmodu. i

```

*----- KCMODU - Range Modulator -----*
*
INTEGER MSTPMD      ! max number of steps
PARAMETER (MSTPMD=64)
*
INTEGER IVOLMD      ! volume number
INTEGER NSTPMD      ! number of steps
REAL THICMD (MSTPMD) ! effective thickness of each step
REAL RATIMD (MSTPMD) ! intensity at each step
*
COMMON/KCMODU/IVOLMD, NSTPMD, THICMD, RATIMD
*
*----- End of KCMODU -----*

```

レンジモジュレーションの仕方を決めるコマンドブロック

最大 MSTPMD ステップの飛程調整

/KCVOLU/でこのレンジモジュレータに対応する物体番号
モジュレーションステップ数
各ステップの飛程シフト量
各ステップでの照射粒子数

KCMODU あるいは UKINIT で定義される変数

2. 2. 8. kcpati. i

```

*----- KCPATI - Patient -----*
*
INTEGER MXYZPT      ! max number of voxels
PARAMETER (MXYZPT=128*128*128)
*
INTEGER IVOLPT      ! volume number
REAL XYZOPT (3)      ! local origin in VL system
REAL RP2VPT (3, 3)   ! rotation matrix local to VL
REAL RV2PPT (3, 3)   ! rotation matrix VL to local
INTEGER NXZPT (3)    ! number of x, y, z divisions
REAL XYZ1PT (3)      ! outer corner of first voxel [cm]

```

計画用 CT 画像をベースにした患者データのためのコマンドブロック

三次元計算ボクセルの最大数

/KCVOLU/で患者に対応する物体番号
患者座標系の原点のビーム座標系 (KCVOLU) での位置
患者座標系からビーム座標系への回転行列
ビーム座標系から患者座標系への回転行列
ボクセル数
ボクセルの角

<pre> REAL DXYZPT (3) ! (dx, dy, dz) of voxel increment [cm] REAL DENSPT (WXYZPT) ! water equivalent density [gH2O/cm^3] REAL DOSEPT (WXYZPT) ! physical dose [MeV/gH2O] * COMMON/KCPAT1/IVOLPT, XYZOPT, RP2VPT, RV2PPT, NXYZPT, XYZ1PT, DXYZPT COMMON/KCPAT2/DENSPT COMMON/KCPAT3/DOSEPT * *----- End of KCPAT1 ----- </pre>	<p>ボクセル幅 各ボクセルの水等価密度 各ボクセル中心での線量値</p> <p>KBPAT1 あるいは UKINIT で定義される変数 同上 UKLAST で処理するための計算結果</p>
--	---

<pre> 2. 2. 9. kpccon. i *----- KPCCON - Physical Constants ----- * * Default Units: * time - sec * length - cm or g/cm^2 or gH2O/cm^2 * angle - radian * charge - positron charge * energy - MeV * momentum - MeV/c * mass - MeV/c^2 or g * density - g/cm^3 * * REAL*8 COULOM ! e to coulomb conversion factor * REAL*8 GRAM ! MeV/c^2 to g conversion factor * REAL*8 JOULE ! MeV to Joule Conversion Factor * REAL*8 C ! speed of light in cm/sec * REAL*8 ALPHA ! fine structure constant * REAL*8 EMASS ! electron mass in MeV/c^2 * REAL*8 ERADI ! classical electron radius in cm * REAL*8 AVOGAD ! Avogadro's Number (number of particles per mol) * REAL*8 AMU ! atomic mass unit in MeV/c^2 * REAL*8 PMASS ! proton mass in MeV/c^2 * * PARAMETER (COULOM=1.602 177 33 D-19) </pre>	<p>物理定数の定義</p> <p>基本単位は 時間は秒、 長さはセンチメートルあるいはグラム毎平方センチメートルあるいは水等価厚 角度はラジアン、 電荷は陽電子電荷量、 エネルギーは MeV、 運動量は MeV/c 質量は MeV/c²(粒子)あるいはグラム(物体) 密度はグラム毎立方センチメートル</p>
--	--

<pre> PARAMETER (GRAM=1. 782 662 70 D-39) PARAMETER (JOULE=1. 602 177 33 D-13) PARAMETER (C=2. 997 924 58 D10) PARAMETER (ALPHA=1. D0/137. 035 989 5 D0) PARAMETER (EMASS=0. 510 999 06 D0) PARAMETER (ERADI=2. 817 940 92 D-13) PARAMETER (AVOGAD=6. 022 136 7 D23) PARAMETER (AMU=931. 494 32 D0) PARAMETER (PMASS=938. 272 31 D0) </pre> <p style="text-align: right;">* *----- End of KOPCON -----*</p>	
---	--

2. 2. 10. kcpenc. i

<pre> *----- KCPENC - Subbeam Parameters -----* * INTEGER IVOLPB ! volume number for dummy volume INTEGER NXYPB (2) ! number of pixels REAL XY1PB (2) ! corner of the first pixel REAL DXYPB (2) ! pixel size * COMMON/KCPENC/IVOLPB, NXYPB, XY1PB, DXYPB * INTEGER MPPB ! number of parameters PARAMETER (MPPB=20) * INTEGER KGHQB ! charge [e] INTEGER KMASPB ! mass [MeV/c^2] INTEGER KXOPB ! x of beam origin [cm] INTEGER KYOPB ! y of beam origin [cm] INTEGER KZOPB ! z of beam origin [cm] INTEGER KDXPB ! x component of beam direction INTEGER KDYPB ! y component of beam direction INTEGER KDZPB ! z component of beam direction INTEGER KDISPB ! distance from origin INTEGER KFLUPB ! fluence (number of particles) </pre>	<p>サブビーム生成のためのコモンブロック</p> <p>生成平面に対する/KCVOLU/での物体番号 生成平面のピクセル数 ピクセルの角の位置 ピクセルの大きさ</p> <p>KBPATI あるいは UKINIT で定義される変数</p> <p>ビームパラメータの最大数</p> <p>電荷(を与えるインデックス番号) 質量 ビーム原点位置</p> <p>ビーム方向</p> <p>原点から現在の輸送点までの距離 粒子数</p>
---	---

<pre> INTEGER KEKIPB ! kinetic energy [MeV] INTEGER KDEPPB ! integrated water equivalent depth INTEGER KX1PB ! integrated radiation length 1 INTEGER KX2PB ! integrated radiation length 2 INTEGER KAAPB ! mean projected scattering angle squared [rad^2] INTEGER KTTPB ! mean projected displacement squared [cm^2] INTEGER KATPB ! covariance [rad cm] INTEGER KRRPB ! range straggling squared [(gH20/cm^2)^2] * PARAMETER (KXOPB=3, KYOPB=4, KZOPB=5, KDXPB=6, KDYPB=7, KDZPB=8 & , KCHGPB=1, KMASPB=2, KFLUPB=9, KEK1PB=10, KDI1PB=11, KAAPB=12 & , KATPB=13, KTTPB=14, KRRPB=15, KDEPPB=16, KX1PB=17, KX2PB=18) * *----- End of KCPENC -----*</pre>	<p>運動エネルギー 積分水等価深さ 積分散乱パラメータ1 積分散乱パラメータ2 横方向の1軸に射影した散乱角の分散 横方向の1軸上の散乱変位の分散 相関項 レンジストラグリングの二乗</p> <p>各ビームパラメータのインデックス番号</p>
--	--

2.2.11. kophys. i

<pre> *----- KOPHYS - Physics options -----* * INTEGER ISCAPH ! scattering ! 1: Highland ! 2: Lynch-Dahl INTEGER ILETPH ! LET calculation ! 1: theoretical step by step calculation ! 2: theoretical depth-dose data ! 3: experimental depth-dose INTEGER ISTRPH ! range straggling ! 0: no additional straggling ! 1: gaussian theory ! 2: 1.1% of range INTEGER INUCPH ! nuclear interaction 0: off, 1: on REAL ECUTPH ! cut off energy [MeV] REAL RELKPH ! max of relative kinetic energy change REAL STMIPH ! minimum step size [cm] REAL STMAPH ! maximum step size [cm] REAL SIGRPH ! range in sigma to define longitudinal spread</pre>	<p>物理モデルのオプション</p> <p>散乱のモデル 1: ハイランドの式 2: リンチ-ダールの式 LET 計算法 1: ステップ毎の理論計算 2: テーブル化した理論計算 3: 実験値テーブル/KODED0/の利用 レンジストラグリング 0: なし(ILETPH=3の時) 1: ガウシアン理論計算 2: 水等価深さの1.1% 核反応の効果(0: なし、1: あり) 陽子のカットオフ運動エネルギー ステップあたり運動エネルギー変化率の最大値 最小ステップサイズ 最大ステップサイズ 線量計算領域を決めるための、標準偏差を単位としたピーク深さからの最大深さ</p>
--	--

```

REAL SIGTPH ! range in sigma to define lateral spread
REAL DSIZPH ! dosimeter size (negative -> voxel size) [cm]
*
COMMON/KCPHYS/ISCAPH, ILOSPH, ISTRPH, INUCPH
&, ECUTPH, RELKPH, STMIPH, STMAPH, SIGRPH, SIGTPH, DSIZPH
*
*----- End of KCPHYS -----*

```

同、横方向の最大半径
線量計算の仮想線量計の大きさ(負ならばボクセルサイズ)

KBPHYS あるいは UKINIT で定義される変数

2.2.12. *kcvo/u. i*

```

*----- KCVOLU - 1-Dimensional Volume Definitions -----*
*
* Coordinate system in this common
*   Origin = isocenter
*   Z direction = -beam direction
*
* INTEGER MAXVVL           ! max number of volumes
* PARAMETER (MAXVVL=20)
*
* REAL RV2GVL (3, 3)      ! rotation matrix from VL to global
* REAL RG2VVL (3, 3)      ! inverse matrix
* INTEGER NVOLVL          ! number of defined volumes
* CHARACTER*4 NAMEVL (MAXVVL) ! volume name
* CHARACTER*4 CTYPVL (MAXVVL) ! volume type
* INTEGER ISUBVL (MAXVVL) ! index in volume-type specific common
* INTEGER IMATVL (MAXVVL) ! material number
* REAL ZPOSVL (MAXVVL)    ! volume center height [cm]
* REAL THICVL (MAXVVL)    ! flat surface height [cm] for bolus
* REAL PHIVL (MAXVVL)     ! volume thickness [cm]
*                           ! rotation angle about z [rad]
*
* COMMON/KCVOLU/RV2GVL, RG2VVL, NVOLVL, NAMEVL
&, CTYPVL, ISUBVL, IMATVL, ZPOSVL, THICVL, PHIVL
*
* Volume type definitions
* CHARACTER*4 CWOBVL      ! wobblers system

```

ビームライン上の物体を定義するためのコモンブロック

ここでの座標系(ビーム座標系)はアイソセンターを原点に、入射ビームの反対の方向を z 方向とする。

定義できる物体の最大数

ビーム座標系から治療室座標系への変換するための回転行列
逆行列

定義された物体の数

各物体の名前

各物体の種類コード

サブインデックス(種類による)

材質番号

物体の中心位置の高さ

ボータスの場合はフラット面の高さ

物体の厚さ

ビーム軸回りの回転角

UKINIT で定義される変数

種類コード名

ワブラー

<pre> CHARACTER*4 CPENVL ! subbeam generation CHARACTER*4 CDEGVL ! degrader CHARACTER*4 CMODVL ! range modulator CHARACTER*4 CCOLVL ! collimator CHARACTER*4 CBOLVL ! bolus CHARACTER*4 CPATVL ! patient PARAMETER (CWOBVL='WOBB', CPENVL='PENC', CDEGVL='DEGR', CMODVL='MODU' &, CCOLVL='COLL', CBOLVL='BOLU', CPATVL='PATI') * *----- End of KCVOLU -----*</pre>	<p>サブビーム生成面 ディグレーダ レンジモジュレータ コリメータ ボラス 患者</p>
--	--

2. 2. 13. kwobb. i

<pre> *----- KWOBBS - Wobbler System -----* * Simplest approximation of circular wobbler system * INTEGER IVOLWB ! volume number for wobbler system REAL ANGLWB ! wobbling angle [rad] * COMMON/KWOBBB/IVOLWB, ANGLWB * *----- End of KWOBBS -----*</pre>	<p>ワブラー条件を与えるコモンブロック 磁石の厚さ、X、Y方向の位置の差を無視 /KCVOLU/でワブラーに対応する物体の番号 曲げる角度 UKINITで定義される変数</p>
---	---

2. 3. ブロックデータ

2. 3. 1. kbbeam. f

<pre> ***** BLOCK DATA KBBEAM ! block data for /KOBEBAM/ * IMPLICIT NONE * * ! 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kcpcon.i' INCLUDE 'kcbeam.i'</pre>	<p>/KCBEBAM/に初期値を与えるブロックデータ</p>
--	---------------------------------

```

DATA CHRGBM /1./
DATA AMASBM /PMASS/
DATA FLUEBM /1./
DATA EKINBM /230./
DATA SEKBM /0./
DATA XYZBM /3*0./
DATA UVWBM /3*0./
END

```

2.3.2. *kbmate.f*

```

*****
BLOCK DATA KBMATE ! block data for /KOMATE/
*
* IMPLICIT NONE
*
* Default material definitions
*
* | 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) *
*****
INCLUDE 'kcmate.i'

DATA IAIRMT /1/
DATA IWATMT /2/
DATA NAMEMT /'AIR
& 'LUCITE
& 'ALUMINUM
& 'COPPER
& 'TUNGSTEN
& '10*'

'WATER
'POLYETHYLEN
'IRON
'BRASS
'LEAD
'/'

DATA ZOVAMT /0.4998, 0.5551, 0.5394, 0.5703, 0.4818,
& 0.4656, 0.4564, 0.4552, 0.4025, 0.3958,
& 10*0./

```

/KOMATE/に初期値を与えるブロックデータ

とりあえずよく使いそうな10個だけ定義されている。

<pre> DATA EXTEMT /85.7E-6, 75.0E-6, 74.0E-6, 57.4E-6, 166.4E-6, & 286.0E-6, 321.9E-6, 333.6E-6, 725.2E-6, 820.0E-6, & 10*0./ DATA DENSMT /0.0012, 1.0000, 1.1900, 0.9400, 2.7000, & 7.8700, 8.9600, 8.4890, 19.300, 11.350, & 10*0./ DATA RADLMT /36.66, 36.08, 40.55, 44.78, 24.01, & 13.84, 12.86, 12.30, 6.76, 6.37, & 10*0./ DATA ABSLMT / 90.0, 84.9, 83.6, 78.8, 106.4, & 131.9, 134.9, 141.70, 185.0, 194.0, & 10*0./ END </pre>	
--	--

2.3.3. *kbmodu.f*

<pre> ***** BLOCK DATA KBMODU ! block data for /KCMODU/ * IMPLICIT NONE * * 1970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** INCLUDE 'komodu.i' DATA IVOLMD/0/ DATA NSTPMD/1/ DATA THICMD/MSTPMD*0./ DATA RATIMD/MSTPMD*1./ END </pre>	<p>/KCMODU/に初期値を与えるブロックデータ</p> <p>レンジモジュレーションなし。</p>
--	--

2. 3. 4. *kcpati.f*

<pre> ***** BLOCK DATA KBPATI ! block data for /KCPATn/ * IMPLICIT NONE * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kcpati.i' DATA IVOLPT/0/ DATA XYZOPT/0,0,0./ DATA RP2VPT/1,0,0,0,1,0,0,0,0,1./ DATA RV2PPT/1,0,0,0,1,0,0,0,0,1./ DATA NXYZPT/3*1/ DATA XYZ1PT/3*0./ DATA DXYZPT/3*1./ DATA DENSPT/MXYZPT*1./ DATA D0SEPT/MXYZPT*0./ END </pre>	<p>/KCPAT1/, /KCPAT2/, /KCPAT3/に初期値を与えるブロックデータ</p>
---	--

2. 3. 5. *kbphys.f*

<pre> ***** BLOCK DATA KBPHYS ! block data for /KBPHYS/ * IMPLICIT NONE * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kphys.i' DATA ISCAPH /2/ DATA ILETPH /1/ </pre>	<p>/KBPHYS/に初期値を与えるブロックデータ</p> <p>散乱は Lynch-Dahl の式 LET は各ビームについて計算</p>
--	---

<pre> DATA ISTRPH /1/ DATA INUJPH /1/ DATA ECUTPH /1.0/ DATA RELKPH /0.1/ DATA STMIPH /0.1/ DATA STMAPH /5.0/ DATA SIGRPH /3.0/ DATA SIGTPH /3.0/ DATA DSIZPH /-1.0/ END </pre>	<p>レンジストラグリングはガウシアン理論 核反応の効果を含む 1MeV まで輸送する 10%までの運動エネルギー変化を許す 最小ステップサイズ 100 ミクロン 最大ステップサイズ 5cm 深さ方向にはビームから3シグマまでとる 横方向にはビーム中心から3シグマとる 線量計のサイズはボクセルサイズ</p>
--	--

2.4. メインプログラム

2.4.1. *kpm* in. f

<pre> ***** PROGRAM KPMAIN ! the main program for wobblers method * IMPLICIT NONE * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kcwobb.i' ! wobblers INCLUDE 'kbeam.i' ! beam INCLUDE 'kvolu.i' ! beamline devices INCLUDE 'kcmate.i' ! material data INCLUDE 'kmodu.i' ! range modulation INCLUDE 'kpcenc.i' ! subbeam generation INCLUDE 'kcoll.i' ! collimator REAL RKSTOP INTEGER IU, IV, IVOL, NVOL, ISTEP REAL PARM1 (MPPB), PARM2 (MPPB) REAL ZPOS, PHI, SPHI, CPHI, U, V, R, FLUX, ANGL, SIGM, BESQ, CHSQ, ZW, AW, DW REAL AA, AT, TT </pre>	<p>メインプログラム</p>
---	-----------------

<pre> LOGICAL LSPR C User setup routine CALL UKINIT C Initialize look-up tables CALL KSINIT C Repeat number of range modulation steps DO ISTEP = 1, NSTPMD C Change range modulator thickness IF (IIVOLMD.NE.0) THEN THICVL(IIVOLMD) = THICMD(ISTP) END IF C beta squared and charge squared BESQ = EKINBM*(EKINBM+2.*AMASBM)/(AMASBM+EKINBM)**2 CHSQ = CHRGBM*CHRGBM C Transport the beam through beamline devices PARM1(KOHPB) = CHRGBM PARM1(KMASPB) = AMASBM PARM1(KXOPB) = XYZBM(1) PARM1(KYOPB) = XYZBM(2) PARM1(KZOPB) = XYZBM(3) PARM1(KDXPB) = UVWBM(1) PARM1(KDYPB) = UVWBM(2) PARM1(KDZPB) = UVWBM(3) PARM1(KFLUPB) = FLUEBM*RATIMD(ISTP) PARM1(KEKIPB) = EKINBM PARM1(KDISPB) = 0. PARM1(KDEPPB) = 0. PARM1(KX1PB) = 0. PARM1(KX2PB) = 0. </pre>	<p>各コモン変数の定義をおこなう</p> <p>各関数のルックアップテーブルを初期化する</p> <p>レンジモジュレーションのステップでループ</p> <p>レンジモジュレータのディグレダとしての厚さを変更</p> <p>β の二乗と電荷の二乗を計算しておく</p> <p>PARM1 は親ビームのビームパラメータを含む配列</p>
--	---

<pre> PARM1 (KAAPB) = 0. PARM1 (KATPB) = 0. PARM1 (KTTPB) = 0. PARM1 (KRRPB) = (SEKBM/RKSTOP (CHSQ, BESQ, ZOVAMT (IWATMT) , EXTEMT (IWATMT)))**2 LSPR = .TRUE. IF (IVOLPB.GT.0) THEN NVOL = IVOLPB ELSE NVOL = NVOLVL END IF DO IVOL = 1, NVOL CALL KSTRAK(LSPR, IVOL, PARM1) END DO </pre>	<p>通常の輸送ではビーム拡大を可能にする</p> <p>サブビーム生成面があればそこまで、なければ最後まで輸送する</p> <p>物体単位でビーム輸送をする</p> <p>サブビーム生成を行う</p> <p>親ビームのパラメータ</p> <p>ワブラーのパラメータ</p>
<pre> C Generate and transport subbeams IF (IVOLPB.GT.0) THEN ZPOS = ZPOSVL (IVOLPB) PHI = PHIVL (IVOLPB) SPHI = SIN (PHI) GPHI = COS (PHI) AA = PARM1 (KAAPB) AT = PARM1 (KATPB) TT = PARM1 (KTTPB) IF (IVOLWB.GT.0) THEN ZW = ZPOSVL (IVOLWB) ! wobbling height AW = ANGLWB ! wobbling angle DW = ABS (ZW-ZPOS) ! wobbling distance ELSE AW = 0. DW = 0. END IF DO IV = 1, NXYPB (2) V = XY1PB (2)+DXYPB (2)*(REAL (IV)-0.5) DO IU = 1, NXYPB (1) </pre>	

<pre> U = XY1PB(1)+DXYPB(1)*(REAL(1U)-0.5) R = SQRT(U*U+V*V) IF (TT.GT.0.) THEN ! wobbling/scattered beam CALL KSWOBB(AW, DW, AA, AT, TT, R, FLUX, ANGL, SIGM) ELSE ! broad parallel beam option FLUX = 1./DXYPB(1)/DXYPB(2)/REAL(NXYPB(1)*NXYPB(2)) ANGL = 0. SIGM = SQRT(AA) END IF PARM2(KCHGPB) = PARM1(KCHGPB) PARM2(KMASPB) = PARM1(KMASPB) PARM2(KX0PB) = CPHI*U-SPHI*V PARM2(KY0PB) = SPHI*U+CPHI*V PARM2(KZ0PB) = ZPOS IF (R.GT.0.) THEN PARM2(KDXPB) = ANGL*PARM2(KX0PB)/R PARM2(KDYPB) = ANGL*PARM2(KY0PB)/R ELSE PARM2(KDXPB) = 0. PARM2(KDYPB) = 0. END IF PARM2(KDZPB) = SIGN(SQRT(1.-ANGL*ANGL), UVWBM(3)) PARM2(KDISPB) = 0. PARM2(KFLUPB) = PARM1(KFLUPB)*FLUX*DXYPB(1)*DXYPB(2) PARM2(KEK1PB) = PARM1(KEK1PB) PARM2(KDEPPB) = PARM1(KDEPPB) PARM2(KX1PB) = PARM1(KX1PB) PARM2(KX2PB) = PARM1(KX2PB) PARM2(KAAPB) = SIGM**2 PARM2(KTTPB) = (DXYPB(1)**2+DXYPB(2)**2)/24. ! minimal IF (TT.GT.0.) THEN PARM2(KATPB) = AT/TT*PARM2(KTTPB) ! preserve FD ELSE PARM2(KATPB) = 0. END IF </pre>	<p>二次元メッシュの各グリッドの位置</p> <p>その点でのフラックス、ビーム方向、広がりを計算</p> <p>TT を 0 にしたとき(最初の物体がサブビーム生成面するとき)、平行ビームを生成する</p> <p>PARM2 はサブビームのビームパラメータを含む配列</p> <p>ANGL からビーム方向ベクトルを計算</p> <p>フラックスにピクセル面積をかけて粒子数を得る</p> <p>積分パラメータは親のものに加算していく</p> <p>ピクセル内に一様分布としてビームの広がりを与える</p> <p>ビームの焦点距離を親ビームと同じにする</p>
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<pre> PARM2(KRRPB) = PARM1(KRRPB) IVOL = IVOLPB DO WHILE (IVOL. LT. NVOLVL .AND. PARM2(KEK1PB)*PARM2(KFLUPB). GT. 0.) IVOL = IVOL+1 IF (IVOLCO. NE. 0. AND. IVOL. LE. IVOLCO) THEN LSPR = .FALSE. ! do not spread before collimator ELSE LSPR = .TRUE. END IF CALL KSTRAK(LSPR, IVOL, PARM2) END DO END DO END DO END IF END DO ! range modulation step 1STP C User analysis routine CALL UKLAST END </pre>	<p>サブビーム生成面から以後の物体にビームを輸送する</p> <p>コリメータまではビーム拡大をしない</p> <p>ビームを輸送するサブルーチンを呼ぶ</p> <p>線量分布の解析や出力を行う</p>
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2.5. サブルーチン

2.5.1. *ksinit.f*

<pre> ***** SUBROUTINE KSINIT ! initialize look-up tables * * IMPLICIT NONE * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** REAL RKERF, RKIERF, RKRESR, RKEKIN, RKLET ! functions REAL R ! dummy variable </pre>	<p>メインプログラムでUKINIT後に呼ばれルックアップテーブルの初期化をさせるルーチン</p>
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<pre> R = RKERF(-1.) ! error function R = RKIERF(-1.) ! integrated error function R = RKRESR(-1.) ! residual range as a function of energy R = RKEKIN(-1.) ! kinetic energy as a function of range R = RKLET(-1.) ! depth-dose table RETURN END </pre>	<p>各関数についてはそれぞれのコードを参照</p>
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<pre> 2. 5. 2. ksmear. f ***** SUBROUTINE KSMEAR(N, XLOW, XSIG, Y, YS) ! smears a plot by sigmas * IMPLICIT NONE INTEGER N ! <- Number of bins REAL XLOW(*) ! <- Lower edge of each bin ! XLOW(N+1) must be given to give upper limit REAL XSIG(*) ! <- Sigma X at each bin REAL Y(*) ! <- Average height of each bin REAL YS(*) ! -> Smearred (averaged) height of each bin * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** INCLUDE 'kconcon. i' REAL RKERF ! error function REAL RKIERF ! integrated error function INTEGER I, J, IDIR REAL SI, XI1, XI2, XI, SJ, XJ1, XJ2, XJ, WJI, DXI, DXJ, S11, S12, S21, S22, DXIJ REAL*8 WI, YI C Loop on bin of interest I DO I = 1, N IF (XLOW(I+1).LT.XLOW(I)) THEN </pre>	<p>可変幅ヒストグラムを各ビンで与えられるシグマに応じてガウシアンでスミアするルーチン</p> <p>ビンの数 各ビンの最小値 ただし、最大ビンの最大値を決めるため XLOW(N+1)が与えられなければならない 各ビンでのシグマ 各ビンの高さ スミアした後の各ビンの高さ</p> <p>YSを計算したいビンに着目する</p>
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<pre> XLOW(I+1) = XLOW(I) END IF X11 = XLOW(I) X12 = XLOW(I+1) X1 = 0.5*(X11+X12) DX1 = X12-X11 SI = M_SQRT2*XSIG(I) C Self contribution (I -> I) IF (SI.LE.0.) THEN WI = DX1 ! delta function ELSE IF (DX1.LT.0.25*SI) THEN ! gaussian at center approximation WI = RKERF(0.5*DX1/SI)*DX1 ELSE WI = SI*RK1ERF(DX1/SI) END IF Y1 = Y(I)*WI C Contributions from neighboring bins (J -> I) DO IDIR = -1, 1, 2 J = I+IDIR WJI = 1. DO WHILE (J.GE.1.AND.J.LE.N.AND.WJI.GT.0.) XJ1 = XLOW(J) XJ2 = XLOW(J+1) XJ = 0.5*(XJ1+XJ2) DXJ = XJ2-XJ1 SJ = M_SQRT2*XSIG(J) DX1J = MIN(ABS(X12-XJ1),ABS(X11-XJ2)) IF (SJ.LE.0.OR.DX1J.GT.2.0*SJ) THEN ! out of range WJI = 0. ELSE IF (DX1J.GT.SJ) THEN ! far bin WJI = DX1*DXJ/M_SQRTPI/SJ*EXP(-(X1-XJ)/SJ)**2) ELSE IF (DXJ.LT.SJ) THEN ! narrow bin WJI = 0.5*DXJ*(RKERF((X12-XJ)/SJ)-RKERF((X11-XJ)/SJ)) </pre>	<p> ビンの最小値 最大値 平均値 ビン幅 ルート2倍のシグマ 自ビンからの寄与 シグマが0ならばデルタ関数 ビン幅が小さいときはビンの中心を平均値とするガウス分布で近似 一般にはビン幅内でフラットに分布する階段関数にガウス分布をコンポジットする ここでWI はビンIの高さ1のコンテンツがあったとしてこれがビンIに残る分の面積 ここでYI はビンIの高さY(I)のコンテンツがあったとしてこれがビンIに残る分の面積 他ビンからの寄与 自ビンに近いものから計算していくための探す方向 遠くにいって、寄与がなくなったらそこで止める ビンJとビンIの最短距離 これが非常に遠いときはもう止める ある程度遠いとき ビンJの中心を平均とするガウス分布で近似しビンIの中心での値で近似 ビンJの幅が小さいとき ビンJの中心を平均とするガウス分布で近似しこれをビンIの範囲で積分 </p>
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<p>一般には、 階段関数をガウス分布でスミアし、これを積分</p> <p>WJ1 はビン J に高さ 1 のコンテンツがあったとしてこれがビン 1 に入る分の面積 YJ1 はビン J に高さ Y(1) のコンテンツがあったとしてこれがビン 1 に入る分の面積 W1, Y1 にそれぞれを加える</p> <p>スミアした高さは Y1 と W1 の比</p>	<pre> ELSE ! near and wide bin S11 = RKIERF ((X11-XJ1)/SJ) S12 = RKIERF ((X12-XJ1)/SJ) S21 = RKIERF ((X11-XJ2)/SJ) S22 = RKIERF ((X12-XJ2)/SJ) WJ1 = SJ*0.5*(S12+S21-S11-S22)*DXJ/DX1 END IF WJ = WJ+WJ1 Y1 = Y1+WJ1*Y(J) J = J+1DIR END DO END DO YS(1) =Y1/W1 END DO RETURN END </pre>
<p>ビームのヒストリー(輸送中の複数点でのビームパラメータのテーブル)を基にこのビームが近傍の計算点に寄与する線量を線量分布に加える</p> <p>ビーム原点の患者座標系での位置 ビーム進行方向の患者座標系での表現 ヒストリーテーブルのデータ数 原点からの幾何学的距離 LET ビームの広がり 通過したボクセル</p>	<pre> ***** SUBROUTINE KSPD0(XYZO, UVW, ND, DISM, SLET, TSIG, KV0X) * * Distribute beam dose in patient (Pol-beam association) * * IMPLICIT NONE REAL XYZO(3) ! beam origin in patient frame REAL UVW(3) ! beam direction in patient frame INTEGER ND ! number of steps or table size REAL DISM(*) ! distances from beam origin REAL SLET(*) ! LET's REAL TSIG(*) ! beam spreads projected to one of lateral axes INTEGER KV0X(*) ! Points of interest on the beam path * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** </pre>

```

INCLUDE 'kcmcon.i'
INCLUDE 'kcpati.i'
INCLUDE 'kcpenc.i'
INCLUDE 'kcpphys.i'

REAL RKINTP
INTEGER IKXYZ

INTEGER MBIN
PARAMETER (MBIN=256)
INTEGER NBIN
REAL WBIN
REAL SLBIN(MBIN)
REAL TTBIN(MBIN)
INTEGER IBIN
REAL BIN

INTEGER ID
INTEGER NBASE(3)
INTEGER IXYZ(3)
REAL DXYZ(3)
INTEGER IAX1, IAX2, IAX3
INTEGER I1, I2, I3
INTEGER IVOX
INTEGER IP1, IP2
INTEGER N2

REAL DISM1
REAL RSQR1
REAL SLET1
REAL TT1
REAL TTMIN
REAL SIGT2

! interpolation function
! voxel number to index function

! max number of bins
! number of bins
! bin width
! SLET table
! TSIG**2 table
! bin number
! (fractional bin number)-0.5

! step number
! base numbers for indices
! central axis voxel indices
! vector to Pol from beam origin
! axis numbers
! voxel indices for IAX1, IAX2, IAX3
! voxel number
! temporary variables for IVOX
! number of voxels at I2

! longitudinal distance to Pol from beam
! radius squared to Pol from beam
! LET at DISM1
! beam spread at DISM1
! practical minimum for TT1
! SIGTPH**2

```

高速計算用一時的なルックアップテーブルの定義

<pre> INTEGER I, I1PREV, I2DIR, I2INI, I2FIN, I3DIR, I3INI, I3FIN REAL X, Y C Base numbers NBASE(1) = 1 NBASE(2) = NXYZPT(1) NBASE(3) = NXYZPT(1)*NXYZPT(2) C Determine the base axis IAX1 = 0 X = 0. DO I = 1, 3 Y = (UVW(1)/DXYZPT(1))**2 IF (Y.GT.X) THEN X = Y IAX1 = I END IF END DO IAX2 = MOD(IAX1, 3)+1 IAX3 = MOD(IAX2, 3)+1 C Minimum beam spread squared IF (DSIZPH.LT.0.) THEN ! limitation due to voxel size TTMIN = (DXYZPT(1)**2*(1.-UVW(1))**2 & +DXYZPT(2)**2*(1.-UVW(2))**2 & +DXYZPT(3)**2*(1.-UVW(3))**2)/24. ELSE ! fixed dosimeter size TTMIN = DSIZPH**2 END IF C Maximum beam spread squared of interest in unit of TT1 SIGT2 = SIGTPH**2 C Prepare lookup table for fast interpolation NBIN = MIN(MBIN, 5*ND) ! 5 times number of steps should be enough </pre>	<p>ボクセル番号とボクセルインデックスの変換用</p> <p>最もインデックスが変化する軸を基本軸とする</p> <p>基本軸でない軸</p> <p>ビームの広がりでの最小値として仮想線量計のサイズあるいはボクセルサイズをとる</p> <p>横方向は SIGTPH シングマ以上は範囲外とするためのカットオフ</p> <p>高速補間のためのルックアップテーブルの準備</p>
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<pre> WBIN = (DISM(ND)-DISM(1))/REAL(NBIN) ID = 1 DO I = 1, NBIN X = DISM(1)+WBIN*(REAL(I)-0.5) ! distance at center of bin DO WHILE (DISM(ID+1).LT.X) ID = ID+1 END DO SLBIN(I) = MAX(0.,RKINTP(1,2,DISM(ID),SLET(ID),X)) TTBIN(I) = MAX(TTMIN,RKINTP(1,2,DISM(ID),TSIG(ID),X)**2) END DO </pre>	<p>LET ビームの広がり の二乗</p> <p>通過したボクセルをみていく</p>
<pre> C Track the beam along the path I1PREV = 0 DO ID = 1, ND IF (KVOX(ID).NE.0) THEN </pre>	<p>基本軸のインデックスが変化したとき</p> <p>基本軸から計算されるボクセル番号 計算点(ボクセル中心)の基本軸の座標</p> <p>通過したボクセルを中心に基本軸に垂直な平面内の各計算点の計算を行う 内側から両方向に計算していき遠くなりすぎたら止める</p>
<pre> C Decompose voxel number into indices DO I = 1, 3 IXYZ(I) = IKPXYZ(KVOX(ID),I) END DO </pre>	
<pre> C Search voxels in a plane at IXYZ(IAX1) I1 = IXYZ(IAX1) IF (I1.NE.I1PREV) THEN I1PREV = I1 IP1 = 1+(I1-1)*NBASE(IAX1) DXYZ(IAX1) = (REAL(I1)-0.5)*DXYZPT(IAX1) +XYZ1PT(IAX1)-XYZ0(IAX1) </pre>	

<pre> I2FIN = 1 END IF DO I2 = I2INI, I2FIN, I2DIR N2 = 0 IP2 = IP1+(I2-1)*NBASE(IAX2) DXYZ(IAX2) = (REAL(I2)-0.5)*DXYZPT(IAX2) +XYZ1PT(IAX2)-XYZ0(IAX2) & DO I3DIR = 1, -1, -2 IF (I3DIR.EQ.1) THEN I3INI = IXYZ(IAX3) I3FIN = NXYZPT(IAX3) ELSE I3INI = IXYZ(IAX3)-1 I3FIN = 1 END IF DO I3 = I3INI, I3FIN, I3DIR IVOX = IP2+(I3-1)*NBASE(IAX3) DXYZ(IAX3) = (REAL(I3)-0.5)*DXYZPT(IAX3) +XYZ1PT(IAX3)-XYZ0(IAX3) & C Distance along beam axis from origin DISM1 = DXYZ(1)*UWV(1)+DXYZ(2)*UWV(2) +DXYZ(3)*UWV(3) IF (DISM1.GT.DISM(ND)) GOTO 11 & C Calculate bin number for the lookup tables BIN = (DISM1-DISM(1))/WBIN+0.5 IBIN = MAX(1,MIN(NBIN-1,INT(BIN+0.5))) C Get TSIG from the look-up table TT1 = TTBIN(IBIN) +(TTBIN(IBIN+1)-TTBIN(IBIN))*(BIN-REAL(IBIN)) & </pre>	<p>第2軸でループ</p> <p>第3軸についても同様</p> <p>第3軸でループ ボクセル番号がきまる 計算点座標もきまる</p> <p>ビーム原点からビーム軸に沿った距離 大きすぎたら次のI2へ</p> <p>この距離に対するルックアップテーブルのビン番号</p> <p>ビームの広がり二乗を補間により求める</p>
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<pre> C Radius squared from beam axis RSQR1 = DXYZ(1)**2+DXYZ(2)**2+DXYZ(3)**2 -DISM1**2 IF (RSQR1.GT.TT1*SIGT2) GOTO 11 Get SLET from the look-up table SLET1 = SLBIN(IBIN) +(SLBIN(IBIN+1)-SLBIN(IBIN))*(BIN-REAL(IBIN)) Add dose (=dE/dZ/rho*d2P/dXdY) from this beam N2 = N2+1 DOSEPT(IVOX) = DOSEPT(IVOX) +SLET1*0.5/M_PI/TT1*EXP(-0.5*RSQR1/TT1) END DO ! I3 CONTINUE 11 END DO ! I3DIR IF (N2.LE.0) GOTO 12 END DO ! I2 CONTINUE 12 END DO ! I2DIR END IF ! change in base axis END IF ! in patient END DO ! end of step ID RETURN END </pre>	<p>ビーム軸からの半径の二乗 大きすぎたら次の I2 へ</p> <p>LET を補間により求める</p> <p>I2 でカウントした計算点数 このビームからの寄与をこの計算点の線量に加算する</p> <p>次の I3 へ I3 をスキップした場合はここに直接来る</p> <p>I2 での計算点がない場合残りの I2 をスキップする 次の I2 へ I2 をスキップしたらここに直接来る</p> <p>次の通過ボクセルへ</p>
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2. 5. 4. ksp tmt. f

<pre> ***** SUBROUTINE KSPTMT(IVOX,ZOVA,EXTE,DENS,RADL,ABSL) * * Get medium parameters in patient * </pre>	<p>患者ボクセルに対する物質データの取得</p>
--	---------------------------

```

*
IMPLICIT NONE
INTEGER IVOX ! <- voxel number
REAL ZOVA ! -> Z/A
REAL EXTE ! -> l
REAL DENS ! -> rho
REAL RADL ! -> X0
REAL ABSL ! -> Lambda l
*
* ! 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) *
*****
INCLUDE 'kcpat.i.'
INCLUDE 'kcmate.i.'
IF (IVOX.EQ.0) THEN ! in the air gap
ZOVA = ZOVAMT(IAIRMT)
EXTE = EXTEMT(IAIRMT)
DENS = DENSMT(IAIRMT)
RADL = RADLMT(IAIRMT)
ABSL = ABSLMT(IAIRMT)
ELSE ! in the compressed water
ZOVA = ZOVAMT(IWATMT)
EXTE = EXTEMT(IWATMT)
DENS = DENSPT(IVOX)
RADL = RADLMT(IWATMT)
ABSL = ABSLMT(IWATMT)
END IF
RETURN
END
*

```

ボクセル番号が 0 のときは空気とする

そうでなければ水とする
ただし、密度は/KCPAT1/の DENSPT で与えられる

2.5.5 kspst.f

```

*****
SUBROUTINE KSPTST(XYZP, UVWP, IVOX, RESP)

```

患者座標系の位置、方向からボクセル番号を計算し、次のボクセルまでの距離を計算する

<pre> IVOX = 0 C In the voxel region, the voxel is the box ELSE DO I = 1, 3 DXYZ(1) = ABS(0.5*DXYZPT(1)) XYZC(1) = XYZIPT(1) + (REAL(1XYZ(1)) - 0.5) * DXYZPT(1) END DO IVOX = IKVVOX(1XYZ(1), 1XYZ(2), 1XYZ(3)) END IF C Translate to the box coordinate DO I = 1, 3 XYZB(1) = XYZP(1) - XYZC(1) IF (IVOX.NE.0.AND.ABS(XYZB(1)).GT.DXYZ(1)) THEN XYZB(1) = SIGN(DXYZ(1), XYZB(1)) END IF END DO C find minimum distance to the box surface RESP = BIG DO I = 1, 3 IF (UVWP(1).NE.0.) THEN IF (IVOX.EQ.0) THEN R = (SIGN(DXYZ(1), -UVWP(1)) - XYZB(1)) / UVWP(1) ! from outside ELSE R = (SIGN(DXYZ(1), UVWP(1)) - XYZB(1)) / UVWP(1) ! from inside END IF IF (R.GE.0..AND.R.LT.RESP) THEN RESP = R END IF END IF END DO RETURN </pre>	<p>ボクセル内なら、このボクセルを箱とする</p> <p>箱の中心を原点とする座標へ移動する</p> <p>数値的な誤差でボクセル外へ出てしまう場合はボクセル内に修正する</p> <p>箱の表面の3面を調べる (他の3面は進行方向の反対となる)</p> <p>ボクセル外なら外側から ボクセル内なら内側から</p> <p>距離の最小となる面を通過する</p>
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END

2. 5. 6. kstrak. f

<pre>***** SUBROUTINE KSTRAK(LSPR, IVOL, P) ! transport a beam * IMPLICIT NONE LOGICAL LSPR ! <- Flag to enable beam spread INTEGER IVOL ! <- next volume number REAL P(*) ! <-> beam parameter array * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kcmcon.i' INCLUDE 'kcphys.i' INCLUDE 'kcvolu.i' INCLUDE 'kcwobb.i' INCLUDE 'kcmate.i' INCLUDE 'kcmodu.i' INCLUDE 'kccoll.i' INCLUDE 'kcbolu.i' INCLUDE 'kcpati.i' INCLUDE 'kcpenc.i'</pre>	<p>ビーム 1 本について 1 つの物体およびその前の空気層を輸送する</p> <p>ボラス・コリメータの順でビーム輸送するときの特別な取り扱い ビームパラメータを含む配列</p>
<pre>C External functions REAL*8 DAAA REAL RKESTR, RKSTOP, RKDFDR, RKENUG, RKLET, RKRESR, RKEKIN</pre>	<p>ビームパラメータのヒストリーを一時的にためておくための配列</p>

<pre> REAL DISM(MD), DEPM(MD), WDEP(MD) ! distance, depth, thickness REAL ULET(MD), SLET(MD) ! unsmearred and smeared LET REAL RSIG(MD), TSIG(MD) ! straggling and spread C Other temporary variables CHARACTER*4 CTYP INTEGER IREG, IMAT, IBOL, IVOX, I, J, IX, IY, IREG2 REAL CHSQ, AMAS, BESQ, RESI, STEP, ZPOS, PHI, THIC, ZOVA, EXTE, CORR &, DENS, RADL, ABSL, XYZ(3), UVW(3), DELT, RESP, DELD, DELW, STOP, SCAT &, WER, CPHI, SPHI, ELOS, ENUC, FABS, DOSE, RESR, TV, EK, FL, EK0, FLO &, XL, YL, XV, YV, AT0, TTO, RRO, PQR(3), PQR0(3), PQRU(3), XLOW(MD) C Cumulative parameters REAL*8 AA, TT, AT, RR, PX1, PX2, DIST, DEPT, AA0, AAC0 C Safety margin for media boundary REAL EPS PARAMETER (EPS=0.0001) FL = P(KFLUPB) EK = P(KEK1PB) IF (EK*FL. LE. 0.) RETURN RESR = RKRESR(EK) ! residual range in g/cm^2 H2O CHSQ = P(KOHGPB)**2 AMAS = P(KMASPB) DIST = P(KDISPB) DEPT = P(KDEPPB) PX1 = P(KX1PB) PX2 = P(KX2PB) AA = P(KAAPB) AT = P(KATPB) TT = P(KTTPB) RR = P(KRRPB) AA00 = AA-DKAA(PX1, PX2) ! correction if AA unequals to DKAA </pre>	<p>次の物体まで確実に輸送するためのステップへの付加量</p> <p>ここからの残飛程</p> <p>配列では使いにくいので普通の変数に入れる</p> <p>角度の分散に散乱以外の効果があるときの補正項</p>
---	--

<pre> CTYP = CTYPVL(IVOL) ! volume type ZPOS = ZPOSVL(IVOL) ! height THIC = THICVL(IVOL) ! thickness PHI = PHIVL(IVOL) ! rotation IMAT = IMATVL(IVOL) ! material DO I = 1, 3 UVW(I) = P(KDXPB+I-1) XYZ(I) = P(KXOPB+I-1)+P(KDISPB)*UVW(I) END DO XV = XYZ(1)+UVW(1)/UVW(3)*(ZPOS-XYZ(3)) ! location in VL coord YV = XYZ(2)+UVW(2)/UVW(3)*(ZPOS-XYZ(3)) TV = SQRT(XV**2+YV**2) CPHI = COS(PHI) SPHI = SIN(PHI) XL = CPHI*XV+SPHI*YV ! location in local coord YL =-SPHI*XV+CPHI*YV IREG2 = 2 IF (CTYP.EQ.CMOBVL) THEN ! wobbler, do nothing ELSE IF (CTYP.EQ.CDEGL) THEN ! degrader, do nothing ELSE IF (CTYP.EQ.CMODVL) THEN ! range modulator, do nothing ELSE IF (CTYP.EQ.CCOLVL) THEN ! collimator IMAT = IAIRMT IX = INT((XL-XY1C0(1))/DXYG0(1)+1.) IY = INT((YL-XY1C0(2))/DXYG0(2)+1.) IF (IX.LT.1.OR.IX.GT.NXYG0(1) & .OR.IY.LT.1.OR.IY.GT.NXYG0(2)) THEN EK = 0. FL = 0. ELSE IF (.NOT.OPENG0(IX,IY)) THEN EK = 0. FL = 0. END IF ELSE IF (CTYP.EQ.CB0LVL) THEN ! bolus </pre>	<p>ビーム座標系(KCVOLU)での位置、方向</p> <p>物体中心の高さでの位置</p> <p>ボースあるいはコリメータのローカル座標系での表現</p> <p>物体の種類に応じた処理</p> <p>コリメータの場合はビームが通過するかどうかを調べ、通過すれば空気とし、しない場合は直ちに止める</p>
---	---

<p>ボーラスの場合はフラット面での位置からピクセルを求め、所定の厚さを与える。</p> <pre> IBOL = ISUBVL(IVOL) IX = INT((XL-XY1B0(1, IBOL))/DXYB0(1, IBOL)+1.) IY = INT((YL-XY1B0(2, IBOL))/DXYB0(2, IBOL)+1.) IF (IX.GE.1.AND. IX.LE.NXYB0(1, IBOL).AND. & IY.GE.1.AND. IY.LE.NXYB0(2, IBOL)) THEN ZPOS = ZPOS+0.5*HIGHB0(IX, IY, IBOL) THIC = ABS(HIGHB0(IX, IY, IBOL)) END IF ELSE IF (CTYP.EQ.CPATVL) THEN ! patient IMAT = IWATMT DO I = 1, 3 PQR0(I) = 0. PQRU(I) = 0. PQR(I) = 0. DO J = 1, 3 PQR0(I) = PQR0(I)+RV2PPT(1, J)*(P(KX0PB+J-1)-XYZOPT(J)) PQRU(I) = PQRU(I)+RV2PPT(1, J)*UVM(J) PQR(I) = PQR(I) +RV2PPT(1, J)*(XYZ(J)-XYZOPT(J)) END DO END DO ELSE IF (CTYP.EQ.OPENVL) THEN ! subbeam generator IREG2 = 1 ! only air gap ELSE WRITE(6,*) 'KSTRAK< Invalid volume: IVOL, CTYP = ', & , IVOL, CTYP IREG2 = 0 ! skip this volume END IF C Region loop (1 for air gap, 2 for the volume) IREG = 1 DO WHILE (IREG.LE. IREG2.AND. EK*FL.GT. 0.) IF (IREG.EQ.1) THEN ! in the air gap before the volume RESI = (ZPOS-XYZ(3)-SIGN(THIC/2., UVM(3)))/UVM(3) ZOVA = ZOVAMT(IAIRMT) EXTE = EXTEMT(IAIRMT) </pre>	<p>患者の場合は患者座標系での原点、位置、方向を求めておく</p> <p>サブビーム生成面は実体なし</p> <p>空気層を輸送し次に本体</p> <p>RESIはこの物体または空気を出るまでの距離</p>
--	--

<pre> DENS = DENSMT (IAIRMT) RADL = RADLMT (IAIRMT) ABSL = ABSLMT (IAIRMT) ELSE ! in the volume RESI = (ZPOS-XYZ (3)+SIGN (THIC/2.,UVW (3)))/UVW (3) ZOVA = ZOVAMT (IMAT) EXTE = EXTEMT (IMAT) DENS = DENSMT (IMAT) RADL = RADLMT (IMAT) ABSL = ABSLMT (IMAT) END IF C Check volume order/overlaps IF (RESI.LT.-2.*EPS) THEN WRITE (6,*) 'KSTRAK> Invalid volume order ' & , IVOL, IREG, ' '//NAMEVL (IVOL)//' ', RESI END IF </pre>	<p>前ステップまでの値を記憶しておく</p> <p>とりあえず物体全体を 1 ステップで動くことを考える</p> <p>患者の場合下部構造としてボクセルがある 境界までの距離を得る ボクセルの物質データを得る</p>
<pre> C Track through the medium RESP = 0. ! distance to next PT voxel ND = 0 ! number of data points in patient DO WHILE (RESI.GT.0.) AAO = AA ATO = AT TTO = TT RRO = RR EKO = EK FLO = FL STEP = RESI+EPS </pre>	<p>患者の場合下部構造としてボクセルがある 境界までの距離を得る ボクセルの物質データを得る</p>
<pre> C For patient, step is limited within voxel IF (CTYP.EQ.CPATVL.AND.IREG.EQ.2) THEN IF (RESP.LE.0.) THEN CALL KSPTST (PQR, PQRU, IVOX, RESP) CALL KSPTMT (IVOX, ZOVA, EXTE, DENS, RADL, ABSL) </pre>	<p>患者の場合下部構造としてボクセルがある 境界までの距離を得る ボクセルの物質データを得る</p>

<pre> END IF STEP = MIN(STEP, RESP+EPS) END IF C Kinetic energy is allowed to be zero in patient IF (FLO*EK0. GT. 0.) THEN C Calculate initial stopping power BESQ = EK0*(EK0+2. *AMAS)/(AMAS+EK0)**2 ! beta squared STOP = RKSTOP(CHSQ, BESQ, ZOVA, EXTE) IF (ZOVA. NE. ZOVAMT (IWATMT) .OR. EXTE. NE. EXTEMT (IWATMT)) THEN WER = STOP/RKSTOP (CHSQ, BESQ, ZOVAMT (IWATMT) .EXTEMT (IWATMT)) ! water equivalent ratio ELSE WER = 1. END IF C Step size under limitations STEP = MIN(STEP, STMAPH , MAX(EPS, STMIPH, EK0/STOP/DENS*RELKPH)) C Angle corrected path length for this step DELT = (1. +AA0)*STEP DELD = DELT*DENS ! path length in g/cm^2 DELW = DELD*WER ! path length in g/cm^2 H2O IF (DELW. GT. RESR-EPS) THEN STEP = STEP*RESR/DELW DELT = DELT*RESR/DELW DELD = DELD*RESR/DELW DELW = RESR END IF C Calculate energy loss and update EK EK = RKEKIN (MAX(0., RESR-DELW)) </pre>	<p>ステップはボクセルの境界までとなる</p> <p>ストラグリングを含めるので患者内では飛程以降の点も計算する必要があるのでない場合</p> <p>ステップ開始時の β 二乗とストップピングパワー</p> <p>g/cm^2 当たりの水等価比</p> <p>水ならば当然 1</p> <p>物体の幾何学的サイズに加え、運動エネルギー変化率からの条件、最大ステップサイズ、最小ステップサイズの条件を課して、ステップサイズを決める</p> <p>ステップに角度分散の効果の補正をしてパスレングスを求める</p> <p>残飛程から平均運動エネルギーを求める</p>
--	---

<pre> CORR = MIN(0.5,0.5*STOP*DELD/EK0) ! correction term CORR = 0 IF (DELTW.LT.1.E-2) THEN ! too small for interpolation ELOS = (1.+CORR)*STOP*DELD ELSE ELOS = EK0-EK END IF IF (EK.LE.ECUTPH.OR.DELW.GE.RESR) THEN ELOS = EK0 EK = 0. END IF C Calculate nuclear interactions and update FL IF (INUCPH.EQ.0.OR.ILETPH.NE.1) THEN FABS = 0. ENUC = 0. ELSE Absorption probability FABS = MIN(1.,RKDFDR(RESR)*DELTW*ABSLMT(IWATMT)/ABSL) ENUC = FABS*RKENUC(EK0) END IF FL = FLO*(1.-FABS) C Calculate unsmearred LET in MeV/g cm^2 in water DOSE = (ELOS+ENUC)*FLO/DELTW*(DELT/STEP) C Calculate scattering and update PX1, PX2, AA, AT, TT, RR IF (ISCAPI.NE.0) THEN PX1 = PX1+DBLE(DELD/RADL*CHSQ *(1.+2.*CORR)/(BESQ*(AMAS+EK0))**2) IF (ISCAPI.EQ.1) THEN PX2 = PX2+DBLE(DELD/RADL) ELSE IF (ISCAPI.EQ.2) THEN PX2 = PX2+DBLE(DELD/RADL*CHSQ*(1.+CORR)/BESQ) END IF </pre>	<p>ステップ内でのエネルギー変化の効果の補正項</p> <p>あまり小さいステップではこの方が精度が良い</p> <p>通常はエネルギーの差から線量への寄与を求める カットオフエネルギー以下になった場合の処理</p> <p>核反応の効果を含めないか既に含まれている場合</p> <p>1つの陽子が吸収される確率 陽子が吸収された場合、落とすエネルギー</p> <p>ストラグリングの効果を含まない水中での LET</p> <p>散乱角計算用のパラメータ 1</p> <p>パラメータ 2(ハイランドの式)</p> <p>パラメータ 2(リンチ=ダールの式)</p>
---	---

<pre> C Update parameters AA, AT, TT, RR AA = DKAA(PX1, PX2) + AAC0 SCAT = (AA-AA0)/DBLE(STEP) IF (LSPR) THEN ! scattering propagation AT = AT+DBLE(AA0*STEP)+DBLE(SCAT/2.*STEP**2) TT = TT+DBLE(2.*AT0*STEP)+DBLE(AA0*STEP**2) +DBLE(SCAT/3.*STEP**3) ELSE ! fan beam effect only TT = TT+DBLE(AT0*STEP) END IF IF (ISTRPH.GT.0.AND.TT0.GT.0.) THEN RR = (SQRT(RR)+M_SQRT2*(AA0-AT0*AT0/TT0)*DELDW)**2 END IF END IF Calculate straggling and update RR IF (ISTRPH.EQ.1) THEN RR = RR+DBLE((WER/(1.+CORR)/STOP)**2 *RKESTR(CHSQ, BESQ, ZOVA)*DELD) ELSE IF (ISTRPH.EQ.2) THEN RR = (SQRT(RR)+DBLE(0.011*DELDW))**2 END IF ELSE ! FLO*EKO C Set step size to 0.5% of depth if EK < 0. STEP = MIN(STEP, MAX(STMIPH, 0.005*DEPT/DENS)) TT = TT+DBLE(AT0*STEP) DOSE = 0. DELT = (1.+AA0)*STEP DELD = DELT*DENS ! path length in g/cm^2 DELDW = DELD*WER ! path length in g/cm^2 H20 END IF C Calculate dose in patient </pre>	<p>散乱角は近似式で得る</p> <p>フェルミアイジス理論による散乱効果の計算</p> <p>散乱効果によるストラグリングの増加(実は無視できるほど小さい)</p> <p>ガウシアン理論によるエネルギー損失のばらつきによるストラグリングの増加</p> <p>水等価深さの 1.1%とする場合</p> <p>飛程以後のビーム輸送</p> <p>飛程の 0.5%ステップとする(ストラグリングの約半分)</p>
---	---

患者領域内では線量を計算するのでヒストリーを残す

```

IF (CTYP.EQ.GPATVL.AND.IVOX.NE.0) THEN
  IF (ND.GE.MD) THEN
    WRITE(6,*) 'KSTRAK> too many data points.'
  ELSE
    ND = ND+1
    KVOX(ND) = IVOX      ! voxel number
    WDEL(ND) = DELW      ! bin width
    DISM(ND) = DIST+DBLE(0.5*STEP) ! mid point in cm
    DEPM(ND) = DEPT+DBLE(0.5*DELW) ! mid point in cmH2O
    ULET(ND) = DOSE      ! unsmeared LET
    IF (ILETPH.NE.1) THEN ! smeared LET
      SLET(ND) = FLO*RKLET(DEPM(ND))*(DELW/STEP)
    END IF
    TSIG(ND) = SQRT(TT0) ! lateral spread in cm
    RSIG(ND) = SQRT(RR0) ! straggling in g/cm^2 H2O
    IF (DOSE.GT.0.) ID = ND ! last energy deposit point
  END IF
END IF ! in patient

C
Proceed one step
DIST = DIST+DBLE(STEP)
RESI = RESI-STEP
RESP = RESP-STEP
DEPT = DEPT+DBLE(DELW)
RESR = RESR-DELW
DO I = 1, 3
  XYZ(I) = XYZ(I)+UVW(I)*STEP
END DO
IF (CTYP.EQ.'PATI') THEN
  DO I = 1, 3
    PQR(I) = PQR(I)+PQRU(I)*STEP
  END DO
END IF

C
Kill this beam when unnecessary

```

LET をルックアップテーブルで計算する場合

ブラッグピーク位置を覚えておく

ビームをステップ分輸送する
 原点からの距離
 物体中の残距離
 ボクセル中の残距離
 積分水等価深さ
 残飛程

ビーム座標系での位置

患者座標系での位置

<p>基本的には運動エネルギーがなくなるか、粒子数がなくなれば直ちにとめるが患者内ではストラグリングの効果等での先まで計算する必要がある</p> <p>ストラグリングを入れない場合は計算される LET が十分小さくなったら止める</p> <p>ストラグリングを入れる場合は SIGRPH で指定される範囲プラス 1 シグマを超えたら止める</p>	<pre> IF (EK*FL. LE. 0.) THEN IF (IREG. EQ. 2. AND. CTYP. EQ. CPATVL. AND. ND. GT. 0) THEN IF (ILETPH. NE. 1) THEN IF (SLET(ND). LT. SLET(1)*0.01) THEN RES1 = 0. END IF ELSE IF (REAL(DEPT). GT. DEPM(ID)+SIGRPH*RSIG(ID)) THEN RES1 = MIN(RESI, RSIG(ID)/DENS) END IF ELSE RES1 = 0. END IF IREG = IREG2 END IF END IF END DO ! WHILE (RES1. GT. 0.) IREG = IREG+1 END DO ! WHILE (IREG. LE. IREG2. AND. EK*FL. GT. 0.) C In patient, calculate dose IF (CTYP. EQ. CPATVL. AND. ND. GT. 0) THEN C Smear depth-dose curve (LET given in MeV/g cm^2 H2O) IF (ILETPH. EQ. 1. AND. ISTRPH. GT. 0) THEN DO I = 1, ND XLOW(I) = DEPM(I)-0.5*WDEP(I) ! low edge END DO CALL KSMEAR(ND-1, XLOW, RSIG, ULET, SLET) END IF C Cut off unnecessary data DO WHILE (SLET(ND). LE. 0.) ND = ND-1 END DO </pre>
<p>物体を抜けるまで繰り返す</p> <p>空気層と物体を計算する</p> <p>患者の場合線量分配をする</p>	
<p>まだ入っていない場合、ストラグリングの効果 LET に入れる</p> <p>ULET をぼかして SLET を計算する</p> <p>ゼロデータ点は必要なし</p>	

<pre> C Distribute dose in patient CALL KSPTD0(PQR0, PQRU, ND, DISM, SLET, TSIG, KVOX) END IF ! CTYP C Return updated parameters P(KFLUPB) = FL P(KEK1PB) = EK P(KDISPB) = DIST P(KDEPPB) = DEPT P(KX1PB) = PX1 P(KX2PB) = PX2 P(KAAPB) = AA P(KATPB) = AT P(KTTPB) = TT P(KRRPB) = RR RETURN END </pre>	<p>ビーム近傍の計算点に線量を付加する</p> <p>アップデートしたビームパラメータを返す</p>
---	---

2.5.7. *kswobb.f*

<pre> ***** SUBROUTINE KSWOBB(AW, DW, AA, AT, TT, R, FLUX, ANGL, SIGM) * * Calculate spreaded beam distribution by Wobbler method * IMPLICIT NONE REAL AW ! <- wobbling angle [rad] REAL DW ! <- wobbling distance [cm] REAL AA ! <- mean projected angle squared [rad^2] REAL AT ! <- mean projected angle*displacement [cm rad] REAL TT ! <- mean projected displacement squared [cm^2] REAL R ! <- distance to beam axis [cm] REAL FLUX ! -> Probability per unit area [cm^-2] </pre>	<p>ワブラー法による拡大ビームの粒子分布を計算するルーチン</p> <p>目的の位置のビーム軸からの距離 親粒子1つあたりで、この位置に単位面積あたりに入る確率</p>
---	---

親ビーム軸から測った子ビームの角度(軸対称のため半径方向となる)
子ビーム角度のガウシアン標準偏差

```

*
* REAL ANGL ! -> average subbeam angle (in r direction) [rad]
* REAL SIGM ! -> sigma of subbeam angle [rad]
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation)
*
*****
INCLUDE 'kcmcon. i'
*****
REAL*8 DKBESI ! modified besel function
REAL*8 BES10 ! modified Bessel function of order 0 at A
REAL*8 DR ! argument for DKBESI
REAL*8 SS ! argument for EXP
REAL RW ! wobbling radius
REAL B1 ! ratio of order 1 divided by order 0
REAL B2 ! ratio of order 2 divided by order 0

RW = AW*DW
DR = R*RW/TT
BES10 = DKBESI (0, DR)
B1 = DKBESI (1, DR) / BES10
B2 = DKBESI (2, DR) / BES10

SS = -0.5 * (R*R+RW*RW) / TT
FLUX = 0.5D0 / (M_PI * TT) * (EXP(SS) * BES10)
ANGL = R*AT / TT - (RW*AT / TT - AW) * B1
SIGM = SQRT (AA - AT*AT / TT + (RW*AT / TT - AW) ** 2 * ((1. + B2) / 2. - B1 * B1))

RETURN
END

```

これらの解析的計算式は外村氏のコード及びレポートに基づいている

2. 6. INTEGER 関数

2. 6. 1. ikpvox. f

```

*****
INTEGER FUNCTION IKPV0X (IX, IY, IZ)

```

ボクセルインデックスからボクセル番号を求める整数関数


```

N = 1
DO I = 1, IXYZ-1
  N = N*XYZPT(1)
END DO
IKPXYZ = MOD((I*VOX-1)/N, N*XYZPT(1*XYZ))+1
RETURN
END

```

2.7. REAL 関数

2.7.1. rkdfdr. f

```

*****
REAL FUNCTION RKDFDR(RESR)
*
* Returns 1/F dF/dR for water, where
* F: fluence
* R: residual range measured in g/cm^2
*
* Ref: Lee and Webb, Phys. Med. Biol. 38(1993)989-998
* where F is parametrized by F(R) = F(0) (1+0.012R)
*
* IMPLICIT NONE
REAL RESR ! residual range in g/cm^2 of water
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation)
*****
RKDFDR = 0.012/(1.+0.012*RESR)
RETURN
END

```

水中での核反応により陽子 1 個、1cm あたり吸収される確率を与える

この論文によると、実験的に、陽子数は深さとともに直線的に減少する。その減少率は飛程での陽子数を 1 とすると 0.012/cm である。

2.7.2. rklet. f

```

*****
REAL FUNCTION RKLET(DEPT)

```

LET に対して実験データあるいは計算したルックアップテーブルを使う場合に参照する LET を与える実数関数

```

*
* LET in water [MeV/g cm^2] as a function of depth
*
*
* IMPLICIT NONE
* REAL DEPT ! depth in g/cm^2 H2O measured from beam source
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation)
*****
INCLUDE 'kcbeam. i'
INCLUDE 'kcmate. i'
INCLUDE 'kcphys. i'
INCLUDE 'kcdedo. i'

REAL RKSTOP, RKDFDR, RKENUC, RKESTR, RKINTP, RKRESR, RKEKIN, RKINT2
INTEGER N
PARAMETER (N=512)
REAL XMIN, Y(N), STEP
DATA STEP/0. /
SAVE Y, XMIN, STEP

REAL E(N), XLOW(N), YO(N), EKIN, RESR, FLUE, CORR, ELOS, BESQ, CHSQ
&, XMID, RSTR, F, ENUC
INTEGER I

C Initialization
IF (STEP.EQ.0.0R.DEPT.LT.0.) THEN
  IF (ILETPH.EQ.3) THEN ! depth-dose given by /KCDED0/
    CORR = RKRESR (EKINBW) - RKRESR (EKINDD)
    XMIN = DEPTDD (1) + CORR
    STEP = (DEPTDD (NDDD) - DEPTDD (1)) / REAL (N)
    DO I = 1, N
      XMID = XMIN + STEP * (REAL (I) - 0.5)
      Y (I) = RKINTP (2, NDDD, DEPTDD, DOSEDD, XMID - CORR)
    END DO
  
```

ルックアップテーブルの定義

実験データを使う場合
 ビームエネルギーが実験と異なる場合の補正
 高速計算のため実験データを補間して等分割点での LET を求める

<pre> ELSE CHSQ = CHRGBM**2 EKIN = EKINBM RESR = RKRESR(EKIN) XMIN = 0. STEP = RESR*1.2/REAL(N) FLUE = 1. RSTR = 0. DO I = 1, N XLOW(I) = XMIN+STEP*REAL(I-1) IF (EKIN.GT.0.) THEN BESQ = EKIN*(EKIN+2.*AMASBM)/(AMASBM+EKIN)**2 RESR = MAX(0., RESR-STEP) ELOS = EKIN-RKEKIN(RESR) IF (EKIN-ELOS.LT.ECUTPH.OR.RESR.EQ.0.) THEN ELOS = EKIN END IF IF (ISTRPH.EQ.1) THEN RSTR = RSTR +RKESTR(CHSQ,BESQ,ZOVAMT(IWATMT))*STEP /RKSTOP(CHSQ,BESQ,ZOVAMT(IWATMT),EXTEMT(IWATMT))**2 E(I) = SQRT(RSTR) ELSE IF (ISTRPH.EQ.2) THEN E(I) = 0.011*(XLOW(I)+0.5*STEP) END IF IF (INUCPH.EQ.1) THEN F = MIN(1.,RKDFDR(RESR)*STEP) ! absorption prob ENUC = RKENUC(EKIN)*F ! absorbed energy ELSE F = 0. ENUC = 0. END IF Y0(I) = (ELOS+ENUC)/STEP*FLUE ! average LET EKIN = EKIN-ELOS FLUE = FLUE*(1.-F) </pre>	<p>理論計算の場合</p> <p>飛程の 1.2 倍まで計算する</p> <p>同様に等分割点での LET をテーブルにする</p> <p>エネルギー損失は残飛程—エネルギー関係から求める</p> <p>エネルギー—ストラグリングの計算</p> <p>核反応の効果</p> <p>LET の計算</p>
---	--

<pre> ELSE Y0(I) = 0. E(I) = SQRT(RSTR) END IF END DO CALL KSMEAR(N-1, XLOW, E, Y0, Y) ! straggling convolution END IF END IF C Lookup table interpolation RKLET = MAX(0., RKINT2(N, XMIN, STEP, Y, DEPT)) RETURN END </pre>	<p>運動エネルギーがゼロでもしばらく計算する</p> <p>ストラグリングでほかして LET を得る</p> <p>高速二次補間</p>
--	---

2.7.3. *rkxde.f*

<pre> ***** REAL FUNCTION RKDXDE(EKIN) * * returns -dX/dE in water [LgH2O/cm^2/MeV] * for Simpson's integration method to calculate range * * IMPLICIT NONE * REAL EKIN ! kinetic energy [MeV] * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kcbbeam.i' INCLUDE 'kcmate.i' REAL RKSTOP, BESQ, CHSQ CHSQ = CHRGBM*CHRGBM </pre>	<p>残飛程-エネルギー関数を求めるためにラウレンツ分布法数値積分のための非積分関数 水中のストップピングパワーの逆数を返す</p>
---	--

<pre> BESQ = EKIN*(EKIN+2.*AMASBM)/(EKIN+AMASBM)**2 RKDXDE = 1./RKSTOP (CHSQ, BESQ, ZOVAMT (IWATMT), EXTEMT (IWATMT)) RETURN END </pre>	
--	--

2.7.4. rkekin.f

<pre> ***** REAL FUNCTION RKEIN (RESR) * * Returns mean kinetic energy at given residual range * * IMPLICIT NONE * REAL RESR ! residual range in water [gH2O/cm^2] * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INCLUDE 'kcbeam. i' INCLUDE 'kcphys. i' REAL RKRESR, RKINT1 INTEGER N, I PARAMETER (N=512) REAL XMIN, STEP, Y (N), EK, RS, EPS, RS1, EK1, EK2 DATA STEP/0./ DATA EPS/1.E-5/ SAVE Y, XMIN, STEP C Initialization IF (STEP.EQ.0..OR.RESR.LT.0.) THEN XMIN = 0. STEP = (SQRT (RKRESR (EKINBM)) - XMIN) / REAL (N) EK1 = ECUTPH </pre>	<p>残飛程から運動エネルギーを求める関数</p> <p>ルックアップテーブルの定義</p> <p>ルックアップテーブルの初期化 残飛程のルートを等分割して対応する運動エネルギーをテーブルにする。</p>
--	--

<pre> DO I = 1, N RS1 = XMIN+STEP*(REAL(I)-0.5) EK2 = EKINBM DO WHILE (ABS(EK1-EK2)/(EK1+EK2).GT.EPS) EK = (EK1+EK2)/2. RS = SQRT(RKRESR(EK)) IF (RS.GT.RS1) THEN EK2 = EK ELSE EK1 = EK END IF END DO Y(I) = EK ! kinetic energy at every SQRT (RESR) bin END DO END IF IF (RESR.LE.0.) THEN RKEKIN = 0. ELSE RKEKIN = MAX(0.,RKINT1(N,XMIN,STEP,Y,SQRT(RESR))) END IF RETURN END </pre>	<p>残飛程 RS1 に対応する運動エネルギーは運動エネルギーから残飛程を求める関数 RKRESR から二分法で求める。</p> <p>残飛程は常に正でゼロなら運動エネルギーもゼロ 残飛程のルートで高速補間</p>
---	---

<p>2.7.5. <i>rkenuc.f</i></p> <pre> ***** REAL FUNCTION RKENUC(EK) * * Returns mean local energy deposition from nuclear interaction * The simplest formulation to reproduce plots in reference * Carlsson et al., Phys. Med. Biol. 42(1997)1033-1053. * IMPLICIT NONE </pre>	<p>1つの陽子の核反応から生ずる荷電二次粒子が落とす全エネルギーを与える</p>
---	---

```

REAL EK ! proton kinetic energy at interaction
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation)
*****
RKENUC = MAX(0., EK-45.)

RETURN
END

```

カーレルソンの論文の彼らのモデル計算による二次粒子からの線量のプロットを再現する最も簡単な近似式。

2. 7. 6. rkerf. f

```

*****
REAL FUNCTION RKERF(X)
*
* Fast look-up-table interpolation for error function ERF(X)
*
* erf(x) = 2/sqrt(pi)*Integral[0->x](exp(-t**2) dt)
* Integral[x, inf](normal(t)/dt) = (1-erf(x/sqrt(2)))/2
* where normal(x) is the normal gaussian distribution function.
*
* IMPLICIT NONE
REAL X
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation)
*****

REAL RKINT1, ERF
INTEGER N, I
PARAMETER (N=256)
REAL Y(N), STEP, XMIN, XMAX, X1
PARAMETER (XMAX=4.)
DATA STEP/0./
SAVE XMIN, Y, STEP

C Initialization

```

誤差関数 ERF に相当する補関関数

関数の機能としては C、FORTRAN の標準拡張関数 erf と同等。プラットフォームによっては単純に erf を使用した方が早い場合が考えられる。その場合は RKERF の代わりに ERF を使用する。

<pre> IF (STEP.EQ.0.) THEN XMIN = 0. STEP = (XMAX-XMIN)/REAL(N) DO I = 1, N X1 = XMIN+STEP*(REAL(I)-0.5) Y(I) = ERF(X1) END DO END IF X1 = ABS(X) IF (X1.LT.0.5*STEP) THEN RKERF = 2.*X*Y(1)/STEP ELSE RKERF = RKINT1(N,XMIN,STEP,Y,X1) IF (X.LT.0.) RKERF = -RKERF END IF RETURN END </pre>	<p>Xを等分割し各点で誤差関数を計算しておく</p> <p>誤差関数は反対称 0で0にするための修正 一次関数で補間</p>
--	---

2.7.7. rkestr. f

<pre> ***** REAL FUNCTION RKESTR(CHSQ, BESQ, ZOVA) ***** * * The lowest order calculation for ionization energy loss * fluctuation squared per unit path length [MeV^2 g^-1 cm^2] * * IMPLICIT NONE REAL CHSQ ! particle charge squared REAL BESQ ! particle beta squared REAL ZOVA ! medium effective Z/A * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** </pre>	<p>エネルギー損失の統計的なふらつきから単位物質長当たりエネルギーストラグリングの二乗を与える関数</p>
---	--

```

INCLUDE 'kcmcon. i'
INCLUDE 'kpccon. i'

REAL FM
PARAMETER (FM=4. *M_P*AVOGAD*ERADJ*ERADJ*EMASS*EMASS)

RKESTR = FM*CHSQ*ZOVA*(1. -0. 5*BESQ)/(1. -BESQ)

RETURN
END

```

ある程度厚い物質ではエネルギーロス分布はガウス分布で近似され、これはその最低次の近似による分散の増加率を与える式

2. 7. 8. rkierf. f

```

*****
REAL FUNCTION RKIERF (X)
*
* Look-up-table interpolation for integrated error function
*
* Ierf (x) = Integrate[0->x] (erf (t)dt)
*
IMPLICIT NONE
REAL X
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation)
*****
REAL RKSIMP, RKERF, RKINT1
EXTERNAL RKERF

INTEGER N, I
PARAMETER (N=256)
REAL Y (N), STEP, XMIN, XMAX, X1, X0
REAL*8 Y1
PARAMETER (XMAX=4.)
DATA STEP/0. /
SAVE XMIN, Y, STEP

```

誤差関数 erf の 0 からの積分値を与える関数

<pre> C Initialization IF (STEP.EQ.0.) THEN XMIN = 0. STEP = (XMAX-XMIN)/REAL(N) X1 = 0. Y1 = 0. DO I = 1, N X0 = X1 X1 = XMIN+STEP*(REAL(I)-0.5) Y1 = Y1+DBLE(RKSI MP(X0, X1, RKIERF)) Y(I) = Y1 END DO END IF X1 = ABS(X) IF (X1.LT.0.5*STEP) THEN RKIERF = 2.*X1*Y(1)/STEP ELSE RKIERF = RKINT1(N, XMIN, STEP, Y, X1) END IF RETURN END </pre>	<p>Xを等分割し、各分割点間は誤差関数をシンプソン積分する。</p> <p>この関数は対称 ゼロでゼロにするための修正 一次関数で補間</p>
---	--

2.7.9. rkint1.f

<pre> ***** REAL FUNCTION RKINT1(N, XMIN, STEP, Y, X) * * Linear interpolation (returns interpolated y at x1) * for equal-sized bin histogram * IMPLICIT NONE INTEGER N ! number of data points REAL XMIN ! lower edge of first bin </pre>	<p>高速一次関数補間を行うための関数</p> <p>補間に使う等ビン幅のヒストグラムのなデータ</p>
--	--

<pre> REAL STEP ! bin size REAL Y(*) ! y at mid point of each bin REAL X ! x point of interest * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** REAL X1 INTEGER I1 X1 = (X-XMIN)/STEP+0.5 ! (fractional bin number)-0.5 I1 = MIN(MAX(INT(X1+0.5), 1), N-1) ! bin number RKINT1 = Y(I1)+(Y(I1+1)-Y(I1))*(X1-REAL(I1)) RETURN END </pre>	<p>高速線形補間計算</p>
---	-----------------

<pre> 2. 7. 10. rkint2. f ***** REAL FUNCTION RKINT2(N, XMIN, STEP, Y, X) * * Parabola interpolation (returns interpolated y at x) * for equal-sized bin histogram * IMPLICIT NONE INTEGER N ! number of data points REAL XMIN ! lower edge of first bin REAL STEP ! bin size REAL Y(*) ! y at mid point of each bin REAL X ! x point of interest * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** REAL XI, DX </pre>	<p>高速二次関数補間を行うための関数</p> <p>補間に使う等ビン幅のヒストグラムのための関数</p>
--	---

<pre> INTEGER I1, I2, I3 X1 = (X-XMIN)/STEP+0.5 ! (fractional bin number)-0.5 I2 = INT(X1+0.5) ! nearest bin number IF (I2.LE.1) THEN RKINT2 = Y(1)+(Y(2)-Y(1))*(X1-1.) ELSE IF (I2.GE.N) THEN RKINT2 = Y(N)+(Y(N)-Y(N-1))*(X1-REAL(N)) ELSE I1 = I2-1 I3 = I2+1 DX = X1-REAL(I2) RKINT2 = Y(I2)+0.5*(Y(I3)-Y(I1))+Y(I3)+Y(I1)-2.*Y(I2))*DX*DX END IF RETURN END </pre>	<p>ヒストグラムの端では線形補間</p> <p>高速パラボラ補間計算</p>
---	---

<pre> 2.7.11. rkintp.f ***** REAL FUNCTION RKINTP(IORD, N, X, Y, X0) * * Linear/parabolic interpolation (returns interpolated y at x0) * * * IMPLICIT NONE INTEGER IORD ! order of interpolation (1:linear or 2:parabolic) INTEGER N ! number of data points ! > 0 X's are in ascending order ! < 0 X's are in descending order REAL X(*) ! x data points REAL Y(*) ! y data points REAL X0 ! x point of interest * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** </pre>	<p>可変間隔のデータ点セットを用いた線形あるいはパラボラ補間計算のための関数</p> <p>N 個の二次元データ点セットからある点での値を補間により求める データ点は X についてソートされている必要がある。</p>
--	---

<pre> INTEGER I1, I2, I REAL X00, X01, X02, Y01, Y02, A, B I1 = 1 I2 = ABS(N) DO WHILE (I2-11. GT. 1) I = (I1+I2)/2 IF (X0*REAL(N). LT. X(I)*REAL(N)) THEN I2 = I ELSE I1 = I END IF END DO IF (IORD.LE. 1) THEN RKINTP = Y(I1)+(Y(I2)-Y(I1))/(X(I2)-X(I1))*(X0-X(I1)) ELSE I = MIN(N-1, MAX(2, I1+INT((X0-X(I1))/(X(I2)-X(I1))+0.5))) X00 = X0-X(I) X01 = X(I-1)-X(I) X02 = X(I+1)-X(I) Y01 = Y(I-1)-Y(I) Y02 = Y(I+1)-Y(I) A = (Y01/X01-Y02/X02)/(X01-X02) B = Y01/X01-A*X01 RKINTP = Y(I)+(A*X00+B)*X00 END IF RETURN END </pre>	<p style="text-align: center;">線形補間の場合</p> <p style="text-align: center;">バラボラ補間の場合</p>
--	---

2. 7. 12. rkresr. f

REAL FUNCTION RKRESR(EKIN)

与えられた運動エネルギーから残飛程を求める補間関数

```

*
* Returns mean residual range [g/cm^2 H2O] at given kinetic energy *
* for particle given in /KCBEAM/ *
* IMPLICIT NONE *
* REAL EKIN ! kinetic energy [MeV] *
*
* 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) *
*****
INCLUDE 'kcbbeam. i'
INCLUDE 'kcphys. i'

REAL RKDXDE, RKSIMP, RKINT1
EXTERNAL RKDXDE

INTEGER N, I
PARAMETER (N=512)
REAL XMIN, STEP, Y(N), EK1, EK2
DATA STEP/0./
SAVE XMIN, STEP, Y
REAL*8 RS

IF (STEP.EQ.0. .OR. EKIN.LT.0.) THEN
  XMIN = ECUTPH
  STEP = (EKINBM-XMIN)/REAL(N)
  RS = 0.
  EK2 = XMIN
  DO I = 1, N
    EK1 = EK2
    EK2 = XMIN+STEP*(REAL(I)-0.5)
    RS = RS+DBLE(RKSIMP(EK1, EK2, RKDXDE))
    Y(I) = RS
  END DO
END IF
END IF

```

補間のためのルックアップテーブル

ルックアップテーブルの初期化

ビームエネルギーからカットオフエネルギーまで等分割

ストップングパワーの逆数を積分して残飛程 0 からエネルギー損失を逆算する

<pre> IF (EKIN.LE.ECUTPH) THEN RKRESR = 0. ELSE RKRESR = MAX(0., RKINT1(N, XMIN, STEP, Y, EKIN)) END IF RETURN END </pre>	<p>カットオフエネルギー以下では残飛程 0 とする</p> <p>通常は高速線形補間</p>
---	---

<pre> 2. 7. 13. rksimp. f ***** REAL FUNCTION RKSIMP(A, B, FUNC) * * Simpson integration of function FUNC between [A, B] * I have no idea how this algorithm works. * I just took it from Dr. Tomura's code. * IMPLICIT NONE REAL A ! lower limit of the integration range REAL B ! upper limit of the integration range REAL FUNC ! real function FUNC(X) to be integrated EXTERNAL FUNC * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) ***** INTEGER MAXI ! maximum 2**MAXI divisions PARAMETER (MAXI=10) REAL E ! max relative error PARAMETER (E=1.E-4) INTEGER INDEX REAL FAFB, XH, XHA, XI, XJ, XNEW, XIP, FNEWX, S, AA, BB </pre>	<p>シンプソン法による数値積分</p> <p>シンプソン法とは任意の関数を二次関数で補間してそれを解析積分する数値積分法 このプログラムは補間するための点の取り方とかを自動的に決めるためのアルゴリズムを含んでおり、積分範囲と被積分関数名を与えるだけで積分値を返す。</p> <p>コード自身は外村氏のコードからとり、アルゴリズムを理解しようとしたが良く分からなかった。</p>
---	---

```

XIP=0.
IF (A. EQ. B) GO TO 1009
AA=MIN (A, B)
BB=MAX (A, B)
FAFB=FUNC (A) +FUNC (B)
XH=BB-AA
XJ=XH*. 5*FAFB
XNEW=AA+XH*. 5
XHA=XH/6.
INDEX=0
1006 FNEWX=FUNC (XNEW)
IF (INDEX. GT. 0) GO TO 1003
INDEX=1
X1=XHA*(FAFB+FNEWX*.4.)
1004 XJ=(XJ+X1*.3.)*. 25
INDEX=INDEX+1
IF (INDEX. GT. MAXI) GO TO 1009
XH=XH*. 5
XNEW=AA+XH*. 5
S=0.
1005 IF (XNEW. LT. BB) GO TO 1006
XIP=(XJ+XH*.2. *S)/3.
IF (ABS (XIP-XI) . LE. ABS (E*XIP)) GO TO 1009
XI=XIP
GO TO 1004
1003 S=S+FNEWX
XNEW=XNEW+XH
GO TO 1005
1009 RKSIMP=XIP
IF (A. GT. B) RKSIMP=-RKSIMP
RETURN
END

```

2. 7. 14. rkstop. f

最低次近似でストッピングパワーを与える関数

<pre> * * REAL FUNCTION RKSTOP(GHSQ, BESQ, ZOVA, EXTE) * * Ionization energy loss for proton per unit path length, * the "stopping power" [MeV g^-1 cm^2] * * Input parameters * * IMPLICIT NONE * REAL CHSQ ! particle charge squared * REAL BESQ ! particle beta squared * REAL ZOVA ! medium effective Z/A * REAL EXTE ! medium effective mean excitation energy in MeV * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** INCLUDE 'kcmcon. i' INCLUDE 'kcpcon. i' REAL FK PARAMETER (FK=4. *M_PI*AVOGAD*ERADI*EMASS) RKSTOP = FK*CHSQ*ZOVA/BESQ &*(LOG(2. *EMASS*BESQ/(1. -BESQ)/EXTE)-BESQ) RETURN END </pre>	<p style="text-align: center;">粒子の性質は電荷二乗と速度二乗</p> <p style="text-align: center;">媒質の性質は原子番号質量数比と平均励起エネルギー</p> <p>陽子に関してはこれより高次の項は十分小さくこれで十分。ICRU のすべての効果を含んだ計算と比べ 250MeV まで飛程では 0. 15mm 以下の差しかない。</p>
---	--

2. 8. REAL*8 関数

2. 8. 1. *dkaa. f*

<pre> ***** REAL*8 FUNCTION DCAA(PX1, PX2) * * Gaussian sigma squared of projected scattering angle * by Highland's or Lynch-Dahl's formula. * </pre>	<p>2 種類の積分物質厚から横方向の 1 軸へ射影した散乱角の分散を与える関数</p>
---	--

<pre> * Reference: Lynch and Dalr, Nucl. Instr. Meth. B58(1991)6-10 * * IMPLICIT NONE REAL*8 PX1 ! Integral (z^2/p^2/beta^2/X0)dX REAL*8 PX2 ! Integral (z^2/beta^2/X0)dX (Lynch-Dahl's formula) ! Integral (1/X0)dX (Highland's formula) * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** REAL*8 S2, EPS DATA S2/13. 6D0/ ! Parameter S2 in Highland's formula [MeV/c] ! which is to give the projection angle's sigma DATA EPS/0. 038D0/ ! Parameter epsilon in Highland's formula DKAA = (S2*(1. D0+EPS*LOG(MAX(1. D-3, PX2))))**2*PX1 RETURN END </pre>	<p>有名なハイランドの近似式とリンチ=ダールの改良版とではログ補正項で使う物質厚の定義の みが異なる。物質厚の計算はこの関数を呼ぶ側で計算する。</p>
--	--

<pre> 2. 8. 2. dkbesi. f ***** REAL*8 FUNCTION DKBESI (K, X) * * Modified Bessel function lk(x) * I have no idea what this is doing. I just took from Dr. Tomura's * BTIS code. * * IMPLICIT NONE INTEGER K ! integer order REAL*8 X ! argument * * 19970725 - N. Kanematsu (Mitsubishi Electric Corporation) * ***** </pre>	<p>改ベッセル関数 lk(x) を与える。ここで k は整数。</p> <p>ベッセル関数の定義や計算法は数学の本を見てください。</p>
--	--

	<pre> INTEGER N, I REAL*8 R, T, W, XX, EPS PARAMETER (EPS=1. D-10) N = ABS(K) R = 1. D0 DO I = 1, N R = R**/DBLE(2*I) END DO XX = 0. 25D0**XX T = 1. D0 W = T I = 0 DO WHILE (ABS(T) .GT. ABS(W*EPS)) I = I+1 T = T**XX/DBLE(I*(I+N)) W = W+T END DO DKBES1 = R*W RETURN END </pre>
--	---

2.9. ユーザールーチン(例)

2.9.1. *ukinit.f*

<pre> SUBROUTINE UKINIT IMPLICIT NONE INCLUDE 'kcbeam. i' INCLUDE 'kcmodu. i' INCLUDE 'kcpati. i' INCLUDE 'kpcenc. i' INCLUDE 'kcoll. i' </pre>	<p>メインプログラムで最初に呼ばれ、照射条件を与えるためのユーザールーチン</p>
--	--

<pre> INCLUDE 'kcvolu.i' INCLUDE 'kwobb.i' INTEGER I1, I2, I3, IV0X REAL X1, X2, X3, R WRITE(6, '(A,\$)') 'UKINIT> Initialization ----' C Initial Beam CHRGBM = 1.0 ! proton charge AMASBM = 938.27231 ! proton mass FLUEBM = 1.0 ! number of protons EKINBM = 150. ! kinetic energy SEKBM = 0. ! energy spread XYZBM(1) = 0. ! incident point XYZBM(2) = 0. ! XYZBM(3) = 320. ! (height) UVWBM(1) = 0. ! incident direction UVWBM(2) = 0. ! UVWBM(3) = -1. ! C Range modulation (none) IVOLMD = 0 NSTPMD = 1 THICMD(1) = 0. RATIMD(1) = 1. C Volume definitions NVOLVL = 0 C Wobbler WOB1 NVOLVL = NVOLVL+1 NAMEVL(NVOLVL) = 'WOB1' ! volume name CTYPVL(NVOLVL) = 'WOB1' ! volume type ISUBVL(NVOLVL) = 0 ! subindex </pre>	<p>ビームの定義</p> <p>ここではアイソセンタを原点にビーム方向を-zにする座標系をとっている</p> <p>飛程変調はなし</p> <p>ワブラー</p>
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<pre> IMATVL (NVOLVL) = 1 ! material ZPOSVL (NVOLVL) = 287. ! height THICVL (NVOLVL) = 61. ! thickness PHIVL (NVOLVL) = 0. ! rotation about z IVOLWB = NVOLVL ANGLWB = 0.025 C Scatterer SCT1 NVOLVL = NVOLVL+1 NAMEVL (NVOLVL) = 'SCT1' ! volume name CTYPVL (NVOLVL) = 'DEGR' ! volume type ISUBVL (NVOLVL) = 0 ! subindex IMATVL (NVOLVL) = 10 ! material ZPOSVL (NVOLVL) = 252. ! height THICVL (NVOLVL) = 0.067 ! thickness PHIVL (NVOLVL) = 0. ! rotation about z C Degradar DEG1 NVOLVL = NVOLVL+1 NAMEVL (NVOLVL) = 'DEG1' ! volume name CTYPVL (NVOLVL) = 'DEGR' ! volume type ISUBVL (NVOLVL) = 0 ! subindex IMATVL (NVOLVL) = 3 ! material ZPOSVL (NVOLVL) = 200. ! height THICVL (NVOLVL) = 2. ! thickness PHIVL (NVOLVL) = 0. ! rotation about z C Pencil beams PEN1 NVOLVL = NVOLVL+1 NAMEVL (NVOLVL) = 'PEN1' ! volume name CTYPVL (NVOLVL) = 'PENC' ! volume type ISUBVL (NVOLVL) = 0 ! subindex IMATVL (NVOLVL) = 1 ! material ZPOSVL (NVOLVL) = 22.5 ! height THICVL (NVOLVL) = 0. ! thickness </pre>	<p>空気 中心高さ 厚さ</p> <p>角度 散乱体</p> <p>鉛 高さ 厚さ</p> <p>ディグレーダ</p> <p>ルサイト 高さ 厚さ</p> <p>サブビーム生成面</p> <p>空気 高さ</p>
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<pre> PHIVL (NVOLVL) = 0. ! rotation about z IVOLPB = NVOLVL NXPB (1) = 64 NXPB (2) = 64 XY1PB (1) = -4.61 XY1PB (2) = -4.61 DXYPB (1) = 4.61*2.0/REAL (NXPB (1)) DXYPB (2) = 4.61*2.0/REAL (NXPB (2)) C Collimator COL1 NVOLVL = NVOLVL+1 NAMEVL (NVOLVL) = 'COL1' ! volume name GTYPL (NVOLVL) = 'COLL' ! volume type ISUBL (NVOLVL) = 0 ! subindex IMATVL (NVOLVL) = 1 ! material ZPOSVL (NVOLVL) = 22.5 ! height THICKVL (NVOLVL) = 0. ! thickness PHIVL (NVOLVL) = 0. ! rotation about z IVOLCO = NVOLVL DO I1 = 1, 2 NXYCO (I1) = NXPB (I1) XY1CO (I1) = XY1PB (I1) DXYCO (I1) = DXYPB (I1) END DO DO I1 = 1, NXYCO (1) X1 = XY1CO (1)+DXYCO (1)*(REAL (I1)-0.5) DO I2 = 1, NXYCO (2) X2 = XY1CO (2)+DXYCO (2)*(REAL (I2)-0.5) R = SQRT (X1*X1+X2*X2) IF (R.LT.4.61) THEN OPENCO (I1, I2) = .TRUE. ELSE OPENCO (I1, I2) = .FALSE. END IF END DO END DO </pre>	<p>物体番号 分割数 端の座標 ピクセル幅 コリメータ 高さ ビーム軸回りの回転角 ピクセルはサブビーム面と同一とする</p> <p>ピクセル中心の位置が半径 4.61cm 以下であればビームを通す</p>
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<pre> END DO C Patient PAT1 NVOLVL = NVOLVL+1 NAMEVL (NVOLVL) = 'PAT1' ! volume name CTYPVL (NVOLVL) = 'PAT1' ! volume type ISUBVL (NVOLVL) = 0 ! subindex IMATVL (NVOLVL) = 2 ! material ZPOSVL (NVOLVL) = 0. ! height THICKVL (NVOLVL) = 25. ! thickness PHIVL (NVOLVL) = 0. ! rotation about z IVOLPT = NVOLVL XYZOPT (1) = 0. XYZOPT (2) = 0. XYZOPT (3) = 2.6 NXYZPT (1) = 100 NXYZPT (2) = 100 NXYZPT (3) = 100 XYZ1PT (1) = -10. XYZ1PT (2) = -10. XYZ1PT (3) = -10. DXYZPT (1) = 0.2 DXYZPT (2) = 0.2 DXYZPT (3) = 0.2 IVOX = 0 DO I3 = 1, NXYZPT (3) X3 = XYZ1PT (3)+DXYZPT (3)*(REAL (I3)-0.5) DO I2 = 1, NXYZPT (2) X2 = XYZ1PT (2)+DXYZPT (2)*(REAL (I2)-0.5) DO I1 = 1, NXYZPT (1) IVOX = IVOX+1 X1 = XYZ1PT (1)+DXYZPT (1)*(REAL (I1)-0.5) R = SQRT (X1**2+X2**2) IF (X3.GE.7.0.AND.X3.LT.9.0.AND.R.GE.2.AND.R.LT.3.) THEN DENSPT (IVOX) = 2. </pre>	<p>患者</p> <p>水</p> <p>アイソセンタから高さが±12.5cm の範囲を/KCVOLU/では患者領域とする</p> <p>ビーム座標系で(0., 0., 2.6)の位置に患者座標系の原点を置く</p> <p>ボクセルは 100×100×100 分割</p> <p>角の座標は患者座標系で(-10., -10., -10.)</p> <p>ボクセルサイズは 2mm 立方</p> <p>各ボクセルに密度を与える</p> <p>ボクセル中心位置が患者座標系で 7cm から 9cm の高さ、半径で 2cm から 3cm に入ったらそのボクセルの密度は 2g/cm³ とする</p>
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<pre> ELSE DENSPT (IVOX) = 1. END IF END DO END DO END DO END DO WRITE (6, ' (A) ') ' done.' RETURN END </pre>	<p>そうであれば 1g/cm³</p>
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<p>2.9.2. uk/last. f</p> <pre> SUBROUTINE UKLAST IMPLICIT NONE INCLUDE 'kcpati. i' INTEGER IKPVOX INTEGER NT, ND PARAMETER (NT=50, ND=100) REAL TMAX, DMAX, TMIN, DMIN PARAMETER (TMIN=0., DMIN=0.) PARAMETER (TMAX=10., DMAX=20.0) REAL DOSE (NT, ND), DERR (NT, ND), PQR (3), D, T, XYZ (3), P, Q, R EQUIVALENCE (P, PQR (1)), (Q, PQR (2)), (R, PQR (3)) INTEGER IP, IQ, IR, IVOX, IT, ID, I, J, NENT (NT, ND) INTEGER NH PARAMETER (NH=40000) REAL H (NH) COMMON/PAWC/H </pre>	<p>メインプログラムで最後に呼ばれ、得られた線量分布を処理するユーザーーチン</p> <p>HB00K のためのメモリー領域の確保</p>
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<pre> WRITE(6, '(A, \$)') 'UKLAST> Analysis ----' CALL FLUSH(6) CALL HLIMIT(NH) CALL HB00K2(1, 'Isodose; Radius (cm); Depth (cm)' &, NT, TMIN, TMAX, ND, DMIN, DMAX, 0.) CALL HB00K2(2, 'Isodose Error; Radius (cm); Depth (cm)' &, NT, TMIN, TMAX, ND, DMIN, DMAX, 0.) DO ID = 1, ND DO IT = 1, NT NENT(IT, ID) = 0 DOSE(IT, ID) = 0. DERR(IT, ID) = 0. END DO END DO DO IR = 1, NXYZPT(3) R = XYZ1PT(3)+DXYZPT(3)*(REAL(IR)-0.5) DO IQ = 1, NXYZPT(2) Q = XYZ1PT(2)+DXYZPT(2)*(REAL(IQ)-0.5) DO IP = 1, NXYZPT(1) P = XYZ1PT(1)+DXYZPT(1)*(REAL(IP)-0.5) IVOX = IKV0X(IP, IQ, IR) DO I = 1, 3 XYZ(I) = XYZ0PT(I) DO J = 1, 3 XYZ(1) = XYZ(1)+RP2VPT(I, J)*PQR(J) END DO END DO D = 12.6-XYZ(3) T = SQRT(XYZ(1)**2+XYZ(2)**2) IT = INT((T-TMIN)/(TMAX-TMIN)*REAL(NT)+1.) </pre>	<p>等線量プロット用の二次元ヒストグラムの定義</p> <p>配列の初期化</p> <p>各ボクセルを見ていく</p> <p>ビーム座標系でのボクセルの位置</p> <p>体表面からの深さ ビーム中心軸からの半径 深さビン番号</p>
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<pre> ID = INT((D-DMIN)/(DMAX-DMIN)*REAL(ND)+1.) IF (IT.GE.1.AND.IT.LE.NT .AND.ID.GE.1.AND.ID.LE.ND) THEN NENT(IT, ID) = NENT(IT, ID)+1 DOSE(IT, ID) = DOSE(IT, ID)+DOSEPT(IVOX) DERR(IT, ID) = DERR(IT, ID)+DOSEPT(IVOX)**2 END IF END DO END DO END DO DO ID = 1, ND D = DMIN+(DMAX-DMIN)/REAL(ND)*(REAL(ID)-0.5) DO IT = 1, NT T = TMIN+(TMAX-TMIN)/REAL(NT)*(REAL(IT)-0.5) IF (NENT(IT, ID).GT.0) THEN DOSE(IT, ID) = DOSE(IT, ID)/REAL(NENT(IT, ID)) DERR(IT, ID) = SQRT(MAX(0., DERR(IT, ID)/REAL(NENT(IT, ID)) -DOSE(IT, ID)**2)) END IF CALL HFILL(1, T, D, DOSE(IT, ID)) CALL HFILL(2, T, D, DERR(IT, ID)) END DO END DO CALL HRPUT(0, 'wobbler.hbook', 'NT') WRITE(6, '(A)') 'done.' CALL FLUSH(6) RETURN END </pre>	<p>半径ビン番号</p> <p>カウンタを加算 線量を加算 線量の二乗を加算</p> <p>深さビン番号でループ 半径ビン番号でループ</p> <p>このビンの平均線量 このビンの線量の標準偏差</p> <p>平均線量をプロットに入れる 標準偏差をプロットに入れる</p> <p>ファイルに保存する</p> <p>終わり</p>
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